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Does Finite Knot Energy Lead To Differentiability?<br>by<br>Simon Blatt<br>Philipp Reiter<br>Report No. 12<br>2006

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# Does Finite Knot Energy Lead To Differentiability? 

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#### Abstract

In this article we raise the question if curves of finite ( $j, p$ )-knot energy introduced by O'Hara are at least pointwise differentiable. If we exclude the highly singular range $(j-2) p \geq 1$ the answer is no for $j p \leq 2$ and yes for $j p>2$. In the first case, which also contains the most prominent example of the MöbiUs energy ( $j=2, p=1$ ) investigated by Freedman, He, and Wang, we construct counterexamples. For $j p>2$ we prove that finite-energy curves have in fact a Hölder continuous tangent with Hölder exponent $\frac{1}{2}(j p-2) /(p+2)$. Thus we obtain a complete picture as to what extent the ( $j, p$ )-energy has self-avoidance and regularizing effects for $(j, p) \in(0, \infty) \times(0, \infty)$. We provide results for both closed and open curves.


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## Introduction

A knot energy is a functional that is bounded from below and self-repulsive, i. e. blows up on sequences of embedded curves converging to a curve with a self-intersection [O'H03, Def. 1.1]. One motivation to study such functionals is to find a "nicer", that is, less entangled shape for a given knot in order to determine its knot type, e. g. by following the negative gradient flow of the knot energy up to a local minimum [ He 00 ]. By claiming self-repulsion one hopes ${ }^{1}$ not to run into the danger of leaving the ambient

[^0]isotopy class during this process. Global minimizers within a prescribed knot class may be regarded as optimal representatives of this knot class exhibiting an "ideal" shape; see the nice illustrations of minimizing curves in various knot classes in [KS98].

The idea of considering energy functionals on knots goes back to S. Fukuhara. He thought of the motion of a non-elastic string with electrons on it lying in a viscous liquid absorbing kinetic energy. In a first paper [Fuk88] he treated the case of knotted polygons.

In 1991 J. O'Hara [O'H91] introduced the knot energy ${ }^{2}$

$$
\begin{equation*}
E(\gamma) \quad:=\iint_{X \times X}\left(\frac{1}{|\gamma(s)-\gamma(t)|^{2}}-\frac{1}{D_{\gamma}(s, t)^{2}}\right)|\dot{\gamma}(s)||\dot{\gamma}(t)| \mathrm{d} s \mathrm{~d} t \tag{1}
\end{equation*}
$$

which may be viewed as a functional on the set $A C_{\mathrm{reg}}\left(X, \mathbb{R}^{3}\right)$ of regular absolutely continuous curves defined on an interval $X \subset \mathbb{R}$ or a circle $X=\mathbb{R} /(\ell \mathbb{Z})$. Here $D_{\gamma}(s, t)$ denotes the distance of $\gamma(s)$ and $\gamma(t)$ on $\gamma$. The factor $|\dot{\gamma}(s)||\dot{\gamma}(t)|$ guarantees the invariance under reparametrization, which allows us to restrict our attention to curves parametrized by arc-length, that are naturally Lipschitz continuous.

In 1994 M. Freedman, Z.-X. He, and Z. Wang proved in their seminal paper [FHW94] the existence and $C^{1,1}$-regularity ${ }^{3}$ of $E$-minimizers in prescribed prime knot classes using the invariance of this particular knot energy under Möbius transformations in $\mathbb{R}^{3}$. Due to that they coined the name MÖBIUS energy. Among many other things they proved that $E$ takes finite values on sufficiently smooth embeddings of $\mathbb{S}^{1}$ [FHW94, Prop. 1.5], [He97, Thm. 1.5.1], and that, on the other hand, any curve with finite Möbius energy has locally a bi-Lipschitz constant arbitrarily close to 1 [FHW94, Cor. 1.3]. This means that the Möbius energy exhibits a "regularizing" effect: Finite energy excludes corner points and even more so cusps on the curve. This fact led to the question if finite energy implies differentiability.

In his (unfortunately unpublished) lecture notes on the Möbius energy [He97, pp. 14 - 19] He constructed an open finite-energy curve (of "spiral" shape) that is not differentiable at a boundary point and asked about the differentiability at interior points $[\mathrm{He} 97$, Problem 1.6.3]. ${ }^{4}$

The answer to this question is contained as a special case in our Main Theorem below, which in fact deals with an entire family of energies, the so-called $(j, p)$-energies

$$
\begin{equation*}
E^{j, p}(\gamma):=\quad \mathscr{L}(\gamma)^{j p-2} \iint_{X \times X}\left(\frac{1}{|\gamma(s)-\gamma(t)|^{j}}-\frac{1}{D_{\gamma}(s, t)^{j}}\right)^{p}|\dot{\gamma}(s)||\dot{\gamma}(t)| \mathrm{d} s \mathrm{~d} t \tag{2}
\end{equation*}
$$

[^1]where $\mathscr{L}(\gamma)$ denotes the length of $\gamma$. These energies were introduced and investigated by O'Hara [O'H92a], [O'H92b], [O'H94], [O'H03]. The Möbius energy (1) corresponds to the case $j=2, p=1$. According to the definition at the beginning the general ( $j, p$ )-energy is a true knot energy if and only if $j p \geq 2$, see [ $\mathrm{O}^{\prime} \mathrm{H} 92 \mathrm{~b}$, Thms. 1.9, 2.3].

In 2003 A. Abrams, J. Cantarella, J. H. G. Fu, M. Ghomi, and R. Howard showed that circles are the unique minimizers of all ( $j, p$ )-energies among closed curves if $p \geq 1$ and $(j-2) p<1$ [ACF ${ }^{+} 03$, Cor. 3]. Their proof also works for $(j-2) p \geq 1$, but in this case circles are only "weak minimizers", since their energy (and hence the energy of all closed curves) is infinite in this case, which is a consequence of Lemma 1 (ii). In the $j p<2$ section we will see that circles are no longer minimizers if $p<1$ and $(j-2) p \geq 1$.
Although knot energies are usually defined only in the context of closed curves, the corresponding functionals are obviously also well-defined for open curves, so we always present statements in terms of open and closed curves.

In the present paper we prove the following
Main Theorem (Differentiability of finite-energy curves). Let j, $p \in(0, \infty)$.
(i) If $j p \leq 2$ and $(j-2) p<1$ there are finite-energy curves that are not differentiable. Furthermore these curves can be chosen of "infimal" energy, i. e. with energy arbitrarily close to the energy of a circle or a line respectively.
(ii) If $j p>2$ all finite-energy curves are of class $C^{1, \alpha / 2}$, where $\alpha=(j p-2) /(p+2)$.

With this result we obtain a complete picture of the regularizing effects of O'Hara's $(j, p)$-energies for $(j, p) \in(0, \infty) \times(0, \infty)$; see Figure 1. The graphs of the three functions $j p=2,(j-2) p=1$, and $p=1$ partition the parameter space $(j, p) \in$ $(0, \infty) \times(0, \infty)$ into several regions of completely different behavior: In the white region the ( $j, p$ )-energy has no regularizing effect, see Part (i) of the Main Theorem. In the grey region finite energy does lead to Hölder continuous first derivatives (Main Theorem, Part (ii)), although in the hatched region above the line $p=1$ one cannot hope to find any closed curve with finite energy ${ }^{5}$. (In fact, we expect the same behavior also in the grey hatched region below $p=1$; at least sufficiently smooth curves except lines can be shown to have infinite energy there too, cf. Lemma 1 (ii). ${ }^{6}$ ) In addition we have indicated in Figure 1 the bizarre effect that squares have always finite energy in the white region, whereas the seemingly ideal shape of the circle leads to infinite energy in the white region between the graphs $(j-2) p=1$ and $j p=2$.

[^2]

Figure 1: Range of $(\boldsymbol{j}, \boldsymbol{p})$
$\alpha=(j p-2) /(p+2)$

The paper is structured as follows. In the first section we briefly discuss the range of high singularity $(j-2) p \geq 1$ using techniques that will also appear later on. By means of $\left[\mathrm{ACF}^{+} 03\right.$, Cor. 3], any closed curve has infinite ( $j, p$ )-energy if in addition $p \geq 1$. Furthermore, if an open or closed curve belongs to $C^{3,1}$ on some open subdomain with non-vanishing curvature, its energy is infinite. Note also that by part (ii) of the Main Theorem the energy of polygons (except lines ${ }^{7}$ ) is infinite in the case $j p>2$.

In the second setting $j p<2$ we observe curves of finite energy that are not differentiable, e. g. squares. We obtain the same fact for the boundary case $j=0, p=\infty$ which corresponds to Gromov's distortion [O'H92b, Ex. 1.3]. Our result from the first section is used to understand the bizarre effect in $\{j p<2,(j-2) p \geq 1\}$ mentioned above.

The third section is devoted to the situation $j p=2$, where the involved construction of a non-differentiable curve of finite energy is carried out. This curve possesses a local bi-Lipschitz constant arbitrarily close to 1 , cf. Corollary 4 . We briefly explain the idea of our proof before going into detail. After constructing an open curve that is not differentiable at one inner point and additionally has an arbitrarily small energy, we "glue" it on a cylinder obtaining a closed curve that is not differentiable at one inner point and whose energy is arbitrarily close to that of a circle, which is known $\left[\mathrm{ACF}^{+} 03\right.$, Cor. 3] to be the unique minimizer for $j \leq 2$. This technique of deriving an energy bound for the "closure" of an open curve applies to a wide range of curves. At the end

[^3]of this section we give a detailed proof of the fact that $E^{j, p}$ is lower semi-continuous under several suitable assumptions, generalizing [FHW94, Lemma 4.2].
In the forth section we deal with $j p>2$. In this case any finite energy curve belongs ${ }^{8}$ to $C^{1, \alpha / 2}$, where $\alpha=(j p-2) /(p+2)$. There are two key ingredients for this proof. The first one is a kind of "quantified" bi-Lipschitz constant that was derived in [O'H94, Prop. 1.6]. The second, Lemma 16, was originally stated by S. Semmes and will allow us to develop a technical tool in Lemma 17, which proves the regularity of curves fulfilling the quantified bi-Lipschitz estimate.

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## $(j-2) p \geq 1$

Lemma 1 (Range of high singularity). Let $\gamma: X \rightarrow \mathbb{R}^{3}$ be an open or closed curve in arc-length parametrization, where $X$ denotes an interval of $\mathbb{R}$ or a circle $\mathbb{R} /(\ell \mathbb{Z})$, and $j, p>0$ with $(j-2) p \geq 1$. Then the conditions
(i) $p \geq 1$ and $X=\mathbb{R} /(\ell \mathbb{Z})$ or
(ii) $\gamma \in C^{3,1}$ on an arbitrary open subdomain of $X$, where additionally $\ddot{\gamma} \neq 0$,
both imply $E^{j, p}(\gamma)=\infty$.
Proof. (i) According to Corollary 3 from [ $\left.\mathrm{ACF}^{+} 03\right]$, circles are strict minimizers for $p \geq 1,(j-2) p<1$ among all closed curves of length $2 \pi$ in arc-length parametrization. From the proof of this result one can derive that this statement also holds for all $p \geq 1$, while the circle is only a "weak minimizer" for $(j-2) p \geq 1$, since its energy (and hence the energy of all closed curves) is infinite in this case, which is a consequence from part (ii). So this gives a rigorous proof of the fact that all closed curves in arc-length parametrization have infinite energy for $(j-2) p \geq 1$.
(ii) We start with the Taylor expansion of $\gamma$; for $s, t \in \mathbb{R}$ we obtain

$$
\gamma(s)-\gamma(t)=\quad \dot{\gamma}(t)(s-t)+\frac{1}{2} \ddot{\gamma}(t)(s-t)^{2}+\frac{1}{6} \dddot{\gamma}(t)(s-t)^{3}+R_{4}(s-t)
$$

where $R_{4}(s-t) \leq \frac{1}{4!}\left\|\gamma^{(4)}\right\|_{L^{\infty}\left(R /(\epsilon z), \mathbb{R}^{3}\right)}|s-t|^{4}$. The arc-length parametrization implies $|\dot{\gamma}| \equiv 1,\langle\dot{\gamma}, \ddot{\gamma}\rangle \equiv 0,\langle\dot{\gamma}, \dddot{\gamma}\rangle \equiv-|\dot{\gamma}|^{2}$, so

$$
\begin{equation*}
\frac{|\gamma(s)-\gamma(t)|^{2}}{|s-t|^{2}} \leq 1-\frac{1}{12}|\ddot{\gamma}(t)|^{2}(s-t)^{2}+C_{1}|s-t|^{3} \tag{3}
\end{equation*}
$$

[^4]for some $C_{1}<\infty$ depending on $\left\|\gamma^{(i)}\right\|_{L^{\infty}}, i=1,2,3,4$, and $\ell$. Since $x \mapsto x^{-j / 2}$ is convex on $(0, \infty)$, we obtain $y^{-j / 2} \geq 1-\frac{j}{2}(y-1)$ for all $y \in(0, \infty)$. This leads to
\[

$$
\begin{aligned}
\frac{1}{|\gamma(s)-\gamma(t)|^{j}}-\frac{1}{|s-t|^{j}} & =\frac{1}{|s-t|^{j}}\left[\left(\frac{|\gamma(s)-\gamma(t)|^{2}}{|s-t|^{2}}\right)^{-j / 2}-1\right] \\
& \geq \frac{j}{2}\left(\frac{1}{12}|\ddot{\gamma}(t)|^{2}-C_{1}|s-t|\right)|s-t|^{2-j}
\end{aligned}
$$
\]

Without loss of generality we may assume that there are $c_{1}, \varepsilon_{0}>0$ such that $|\ddot{\gamma}(t)| \geq c_{1}$ for all $t \in\left[0, \varepsilon_{0}\right]$. Now $\varepsilon:=\min \left(\varepsilon_{0}, \frac{c_{1}^{2}}{24 C_{1}}, \frac{1}{2} \ell\right)$ yields
$E^{j, p}(\gamma) \geq \int_{0}^{\varepsilon} \int_{t}^{\varepsilon+t}\left(\frac{1}{|\gamma(s)-\gamma(t)|^{j}}-\frac{1}{|s-t|^{j}}\right)^{p} \mathrm{~d} s \mathrm{~d} t \geq\left(\frac{j}{48} c_{1}^{2}\right)^{p} \varepsilon \int_{0}^{\varepsilon} u^{(2-j) p} \mathrm{~d} u$,
which is infinite if $(j-2) p \geq 1$.

## $j p<2$

The $(j, p)$-energy of the unit square $Q:[0,4] \rightarrow \mathbb{R}^{2}, t \mapsto(t, 0)$ for $t \in[0,1], t \mapsto$ $(1, t-1)$ for $t \in[1,2], t \mapsto(3-t, 1)$ for $t \in[2,3], t \mapsto(0,4-t)$ for $t \in[3,4]$, is finite. The only interesting point is the interaction of neighboring segments, which leads to

$$
\begin{aligned}
& \iint_{[0,1] \times[1,2]}\left(\frac{1}{|Q(s)-Q(t)|^{j}}-\frac{1}{|s-t|^{j}}\right)^{p} \mathrm{~d} s \mathrm{~d} t \leq \iint_{[0,1] \times[1,2]} \frac{\mathrm{d} s \mathrm{~d} t}{|Q(s)-Q(t)|^{j p}} \\
& \quad=\iint_{[0,1]^{2}} \frac{\mathrm{~d} s \mathrm{~d} t}{\left(s^{2}+t^{2}\right)^{j p / 2}} \leq \int_{0}^{\pi / 2} \int_{0}^{\sqrt{2}} \frac{r \mathrm{~d} r \mathrm{~d} \varphi}{\left(r^{2}\right)^{j p / 2}}=\frac{\frac{\pi}{2}}{2-j p} \cdot 2^{(2-j p) / 2}
\end{aligned}
$$

So there is a curve of finite $(j, p)$-energy which is not differentiable.
The same calculation yields that $E^{j, p}$ is not self-repulsive, so this case does not model a knot energy, cf. [O'H92b, Thm. 1.9].
To obtain an open curve with arbitrary small energy take two lines and join them by a sufficiently large angle, i. e. $\pi-\varepsilon$. For closed curves one may replace a small piece of a circle by an angle and adapt the arguments given in the proof of Proposition 13. If $j \in(0,4), p \in\left[1, \frac{2}{j}\right)$ the $(j, p)$-energy of the curves constructed in the next section are bounded by means of HöLDER's inequality in terms of their ( $j, \frac{2}{j}$ )-energy which can be chosen arbitrarily small.
For the boundary case $j=0, p=\infty$ which corresponds to Gromov's distortion [O'H92b, Ex. 1.3] we furthermore obtain

$$
\operatorname{distort}(Q)=\sup _{0 \leq s<t \leq 4} \frac{D_{Q}(s, t)}{|Q(s)-Q(t)|}=2
$$

so that $Q$ is also an example for a non-differentiable curve whose distortion is finite. In the set $\left\{(j, p) \in \mathbb{R}^{2} \mid j>4, \frac{1}{j-2}<p<\frac{2}{j}\right\}$ we obtain the bizarre situation that the energy of a square is finite while the energy of a circle is infinite according to Lemma 1 (ii), so we cannot expect that equation (3) from $\left[\mathrm{ACF}^{+} 03\right]$ also holds for all $p<1$, i. e. that circles are still minimizers of the $(j, p)$-functional.

## $j p=2$

Let $E^{[j]}:=E^{j, 2 / j}$. Note that this notation is (for $j \neq 2$ ) different from $E^{(j)}:=E^{j, 1}$ that is used by O'Hara in [O'H03].

Theorem 2 (Finite $\left(\boldsymbol{j}, \frac{\mathbf{2}}{\boldsymbol{j}}\right.$ )-energy does not imply differentiability). For any $\varepsilon>0$ and $j \in(0,4)$ there is an open curve $\mathbb{R} \rightarrow \mathbb{R}^{2}$ of $\left(j, \frac{2}{j}\right)$-energy $\leq \varepsilon$ parametrized by arc-length that is not differentiable at 0 and coincides outside a neighborhood of $0_{\mathbb{R}^{2}}$ with the $x_{1}$-axis.

Corollary 3 (Version for closed curves). Let $\zeta: \mathbb{R} / \mathbb{Z} \rightarrow \mathbb{R}^{2}$ denote the circle of length 1 . For any $\varepsilon>0$ and $j \in(0,4)$ there is a closed curve $\mathbb{R} / \mathbb{Z} \rightarrow \mathbb{R}^{3}$ in arc-length parametrization whose $\left(j, \frac{2}{j}\right)$-energy lies in the $\varepsilon$-neighborhood of $E^{j, 2 / j}(\zeta)$, but which is not differentiable at 0 .

Using the fact that finite Möbius energy curves possess a local bi-Lipschitz property [FHW94, Cor. 1.3] we immediately deduce

Corollary 4 (Arbitrarily "small" local bi-Lipschitz constant does not imply differentiability). There is an open curve $\mathbb{R} \rightarrow \mathbb{R}^{2}$ in arc-length parametrization that is not differentiable at 0 , but satisfies the following condition:

For any $\varepsilon_{0}>0$ there is a $\delta>0$ such that for any $y \in \mathbb{R}$ the restriction to $[y-\delta, y+\delta]$ is bi-LIPschitz continuous with constant $1+\varepsilon_{0}$.

There is also a closed curve $\mathbb{S}^{1} \rightarrow \mathbb{R}^{2}$ with a similar property.
Before presenting the rigorous proof in several steps we give a brief
Outline of the proof. The main idea is to construct a basic component $\kappa_{\alpha}$, that lies on the $x$-axis outside a neighborhood of the origin.


In a smaller neighborhood of the origin $\kappa_{\alpha}$ lies on a line segment that meets the $x$-axis at the origin in an angle of $\alpha$ as drawn on the left. The energy of $\kappa_{\alpha}$ amounts to $O\left(\alpha^{4 / j}\right)$ by Proposition 7 .

Now let $\alpha^{*}<\alpha$ and replace (a part of) the line segment by a copy of $\kappa_{\alpha^{*}}$ scaled down and restricted to a suitable neighborhood of the origin, so it fits into the gap. Because of the scaling invariance we can choose this copy arbitrarily small, i. e. we change the original curve $\kappa_{\alpha}$ only on a very small subdomain at the origin.


The composite curve meets the origin in an angle of $\alpha+\alpha^{*}$. It turns out that its energy can be estimated in terms of $E^{[j]}\left(\kappa_{\alpha}\right), E^{[j]}\left(\kappa_{\alpha^{*}}\right)$ and some other terms that depend on quantities which can be controlled by the scaling parameter of $\kappa_{\alpha^{*}}$, cf. Proposition 8. So we obtain essentially $O\left(\alpha^{4 / j}\right)+O\left(\left(\alpha^{*}\right)^{4 / j}\right)$ as energy of the composite curve.

By repeating this process inductively for $\alpha_{k}:=\frac{1}{k}$ we obtain a limit curve that performs a rotation of $\sum_{k \in \mathbb{N}} k^{-1}=\infty$ near 0 but has energy $\sum_{k \in \mathbb{N}} k^{-4 / j}<\infty$.

For the rigorous proof of Theorem 2 we will collect a few tools at first which will be used later on, starting with the following fact.

Lemma 5. Let $j>0$. There is a constant $C_{2}$ depending only on $j>0$ such that for $\lambda, \mu>0$

$$
\begin{equation*}
\int_{-\infty}^{-\lambda} \int_{\lambda}^{\infty}\left(\frac{1}{(s-t)^{j}}-\frac{1}{(s-t+\mu)^{j}}\right)^{2 / j} \mathrm{~d} s \mathrm{~d} t \leq C_{2}\left(\frac{\mu}{2 \lambda}\right)^{2 / j} \tag{4}
\end{equation*}
$$

holds.

Proof. Let $q(j, x):=(1-j x)(1+x)^{j}$. Since we have

$$
\frac{\mathrm{d}}{\mathrm{~d} j} q(j, x)=(1+x)^{j}[-x+(1-j x) \log (1+x)]<0 \quad \text { for all } j, x>0
$$

and $j \mapsto q(j, x)$ is continuous on $[0, \infty)$, we arrive at $q(j, x) \leq q(0, x)=1$, i. e.

$$
\begin{equation*}
(1-j x)(1+x)^{j} \leq 1 \quad \text { for all } j, x>0 . \tag{5}
\end{equation*}
$$

For $s, t>0$ and $x=\frac{\mu}{s+t}$ we obtain

$$
\begin{align*}
1-j \frac{\mu}{s+t} & \leq \frac{1}{\left(1+\frac{\mu}{s+t}\right)^{j}}=\left(\frac{s+t}{s+t+\mu}\right)^{j} \\
\Longrightarrow 1-\left(\frac{s+t}{s+t+\mu}\right)^{j} & \leq j \frac{\mu}{s+t} . \tag{6}
\end{align*}
$$

This yields

$$
\iint_{[\lambda, \infty)^{2}}\left(\frac{1}{(s+t)^{j}}-\frac{1}{(s+t+\mu)^{j}}\right)^{2 / j} \mathrm{~d} s \mathrm{~d} t
$$

$$
\begin{aligned}
& =\iint_{[\lambda, \infty)^{2}} \frac{1}{(s+t)^{2}}\left(1-\left(\frac{s+t}{s+t+\mu}\right)^{j}\right)^{2 / j} \mathrm{~d} s \mathrm{~d} t \\
& \stackrel{(6)}{\leq}(j \mu)^{2 / j} \iint_{[\lambda, \infty)^{2}}(s+t)^{-2-2 / j} \mathrm{~d} s \mathrm{~d} t=\frac{1}{\frac{2}{j}} \cdot \frac{1}{1+\frac{2}{j}} \cdot j^{2 / j}\left(\frac{\mu}{2 \lambda}\right)^{2 / j} .
\end{aligned}
$$

By TAYLOR approximation there is a $c_{3}>0$, such that, for any $\alpha \in\left(0, c_{3}\right]$,

$$
\begin{array}{llr}
0<\alpha-\sin \alpha \leq \frac{1}{3} \alpha^{3}, & \frac{1}{2} \alpha \leq \sin \alpha \leq \alpha,  \tag{7}\\
0<1-\cos \alpha \leq \alpha^{2}, & \alpha \leq \tan \alpha \leq 2 \alpha .
\end{array}
$$

Let $\alpha \in\left(0, c_{3}\right]$ and define

$$
\begin{array}{rlrl}
x(\alpha):=\frac{1-\cos \alpha}{\tan \alpha} & <\quad y(\alpha):=\frac{1-\cos \alpha}{\sin \alpha}=\tan \frac{\alpha}{2} & \leq \alpha,  \tag{8}\\
\xi_{\alpha} & :=3 \sin \alpha+x(\alpha) & <\quad \eta_{\alpha}:=3 \alpha+y(\alpha) & \leq 4 \alpha .
\end{array}
$$

This yields by (7)

$$
\begin{align*}
0 \leq y(\alpha)-x(\alpha)=(1-\cos \alpha) \frac{\tan \alpha-\sin \alpha}{\sin \alpha \tan \alpha} & \leq 2 \alpha^{3},  \tag{9}\\
0 \leq \eta_{\alpha}-\xi_{\alpha} & \leq 3 \alpha^{3} .
\end{align*}
$$

Now we are going to construct the components we will later insert into one another inductively.


The basic component $\kappa_{a}: \mathbb{R} \rightarrow \mathbb{R}^{2}$ is defined as drawn above and characterized by the
following explicit formula in arc-length parametrization.

$$
t \mapsto \begin{cases}\binom{\cos \alpha}{\sin \alpha} t & \text { for } t \in[0, y(\alpha)] \\ \binom{\sin (t-\alpha-y(\alpha))}{\cos (t-\alpha-y(\alpha))}+\binom{x(\alpha)+\sin \alpha}{1-2 \cos \alpha} & \text { for } t \in[y(\alpha), y(\alpha)+2 \alpha] \\ \binom{\sin \left(t-\eta_{\alpha}\right)}{-\cos \left(t-\eta_{\alpha}\right)}+\binom{\xi_{\alpha}}{1} & \text { for } t \in\left[y(\alpha)+2 \alpha, \eta_{\alpha}\right] \\ \binom{t}{0}-\binom{\eta_{\alpha}-\xi_{\alpha}}{0} & \text { for } t \in\left[\eta_{\alpha}, \infty\right) \\ -\kappa_{\alpha}(-t) & \text { for } t \in(-\infty, 0)\end{cases}
$$

To compute $E^{[j]}\left(\kappa_{\alpha}\right)$ we need an estimate for the energy of an arc joined with a half line. This curve $\bar{\zeta}_{\alpha}$ was first introduced by Z.-X. He in [He97, Ex. 1.1.2 and p. 15]; his calculations lead to $E^{[2]}\left(\bar{\zeta}_{\alpha}\right)=2-\alpha \cot \frac{\alpha}{2}=O\left(\alpha^{2}\right)$.

## Lemma 6 (Energy of a "hockey stick").

There are constants $c_{4} \in(0,2 \pi), C_{4}<\infty$, depending only on $j>0$, such that for $\alpha \in\left(0, c_{4}\right]$ the curve $\bar{\zeta}_{\alpha}:(-\infty, \alpha] \rightarrow$ $\mathbb{R}^{2}$ given by

$$
t \mapsto\left\{\begin{array}{lll}
\binom{t}{0} & \text { for } & t \leq 0 \\
\binom{\sin t}{1-\cos t} & \text { for } & t \in[0, \alpha]
\end{array}\right.
$$

satisfies

$$
E^{[j]}\left(\bar{\zeta}_{\alpha}\right) \leq C_{4} \alpha^{4 / j} .
$$

Proof. By modifying the arguments from Lemma 1 (ii) (estimate towards the other direction) and applying them to the circle segment we will show that there are constants $c_{5}>0, C_{5}<\infty$ such that $E^{[j]}\left(\left.\bar{\zeta}_{\alpha}\right|_{[0, \alpha]}\right) \leq C_{5} \alpha^{4 / j}$ for any $\alpha \in\left(0, c_{5}\right]$. The $\operatorname{arc} \zeta:=\left.\bar{\zeta}_{\alpha}\right|_{[0, \alpha]}$ obviously fulfills the assumptions of Lemma 1 (ii), so we arrive at

$$
\begin{equation*}
\frac{|\zeta(s)-\zeta(t)|^{2}}{|s-t|^{2}} \geq 1-\frac{1}{12}|\ddot{\zeta}(t)|^{2}(s-t)^{2}-C_{1}|s-t|^{3} \tag{*}
\end{equation*}
$$

for $C_{1}<\infty$ depending only on $\alpha$, since $\left\|\zeta^{(i)}\right\|_{L^{\infty}}=1$ for all $i \in \mathbb{N}$. Since $(1+x)^{-j / 2}=$ $1-\frac{j}{2} \int_{0}^{x}(1+\xi)^{-j / 2-1} \mathrm{~d} \xi$ we obtain

$$
y^{-j / 2} \leq 1-\frac{j}{2} y^{-j / 2-1}(y-1) \leq 1+\frac{j}{4}(1-y) \quad \text { for all } y \in\left[2^{-1 /(j / 2+1)}, 1\right]
$$

This leads to

$$
\begin{aligned}
& \frac{1}{|\zeta(s)-\zeta(t)|^{j}}-\frac{1}{|s-t|^{j}}=\frac{1}{|s-t|^{j}}\left[\left(\frac{|\zeta(s)-\zeta(t)|^{2}}{|s-t|^{2}}\right)^{-j / 2}-1\right] \\
& \leq \frac{j}{4}\left(\frac{1}{12}|\ddot{\zeta}(t)|^{2}+C_{1}|s-t|\right)|s-t|^{2-j} \leq \frac{j}{4}\left(\frac{1}{12}+C_{1} \alpha\right)|s-t|^{2-j}
\end{aligned}
$$

provided $\frac{\mid \zeta\left(s s-\left.\zeta(t)\right|^{2}\right.}{|s-t|^{2}} \geq 2^{-1 /(j / 2+1)}$. But since $\zeta$ has a (uniform) local bi-LIPschitz constant arbitrarily close to 1 , there is a $c_{5}>0$ such that this requirement holds for all $|s-t| \leq$ $c_{5}$. Now

$$
\begin{aligned}
E^{j, p}(\gamma) & \leq \int_{0}^{\alpha} \int_{0}^{a}\left(\frac{1}{|\zeta(s)-\zeta(t)|^{j}}-\frac{1}{|s-t|^{j}}\right)^{p} \mathrm{~d} s \mathrm{~d} t \\
& \leq\left[\frac{j}{4}\left(\frac{1}{12}+C_{1} \alpha\right)\right]^{p} \cdot 2 \int_{0}^{\alpha} \int_{0}^{\alpha-t} u^{(2-j) p} \mathrm{~d} u \\
& \leq \frac{\left[\frac{j}{4}\left(\frac{1}{12}+C_{1} \alpha\right)\right]^{p}}{[(2-j) p+1][(2-j) p+2]} \cdot \alpha^{(2-j) p+2} \stackrel{j p=2}{=} \stackrel{:^{2}}{ } \quad C_{5} \alpha^{4 / j} .
\end{aligned}
$$

Since

$$
\left|\bar{\zeta}_{\alpha}(s)-\bar{\zeta}_{\alpha}(t)\right| \leq\left|\bar{\zeta}_{2 \alpha}(s+\alpha)-\bar{\zeta}_{2 \alpha}(t+\alpha)\right| \quad \text { for } s, t \in[-\alpha, \alpha]
$$

we even obtain $E^{[j]}\left(\left.\bar{\zeta}_{\alpha}\right|_{[-\alpha, \alpha]}\right) \leq C_{5}(2 \alpha)^{4 / j}$ for any $\alpha \leq \frac{1}{2} c_{5}$. Certainly, $E^{[j]}\left(\left.\bar{\zeta}_{\alpha}\right|_{(-\infty, 0]}\right)=0$. It remains to study the interaction of the intervals $(-\infty,-\alpha]$ and $[0, \alpha]$. Let $s \in[0, \alpha]$, $t \in[\alpha, \infty), \alpha \leq c_{4}:=\min \left(c_{3}, \frac{1}{2} c_{5}, 1,(2 j)^{-1 / 2}\right)$, which implies

$$
\begin{equation*}
(1+\tau)^{j} \quad \stackrel{(5)}{\leq} \quad \frac{1}{1-j \tau} \leq 1+2 j \tau \quad \text { for any } \tau \in\left[0, c_{4}^{2}\right] . \tag{10}
\end{equation*}
$$

Since

$$
\left(\frac{(s+t)^{2}}{(\sin s+t)^{2}+(1-\cos s)^{2}}\right)^{1 / 2} \stackrel{(7)}{\leq} \frac{s+t}{s-s^{3}+t} \leq 1+\frac{s^{3}}{t}
$$

and $0 \leq \frac{s^{3}}{t} \leq \alpha^{2} \leq c_{4}^{2}$ we arrive at

$$
\begin{aligned}
& \int_{-\infty}^{-\alpha} \int_{0}^{\alpha}\left(\frac{1}{\left|\bar{\zeta}_{\alpha}(s)-\bar{\zeta}_{\alpha}(t)\right|^{j}}-\frac{1}{(s-t)^{j}}\right)^{2 / j} \mathrm{~d} s \mathrm{~d} t \\
& \quad \leq \int_{\alpha}^{\infty} \int_{0}^{\alpha} \frac{1}{(s+t)^{2}}\left[\left(\frac{(s+t)^{2}}{(\sin s+t)^{2}+(1-\cos s)^{2}}\right)^{j / 2}-1\right]^{2 / j} \mathrm{~d} s \mathrm{~d} t \\
& \stackrel{(10)}{\leq}(2 j)^{2 / j} \int_{\alpha}^{\infty} \int_{0}^{\alpha} \frac{1}{(s+t)^{2}}\left(\frac{s^{3}}{t}\right)^{2 / j} \mathrm{~d} s \mathrm{~d} t \\
& \quad \leq(2 j)^{2 / j} \int_{\alpha}^{\infty} \int_{0}^{\alpha} \frac{s^{6 / j}}{t^{2 / j+2}} \mathrm{~d} s \mathrm{~d} t=\frac{(2 j)^{2 / j}}{\left(\frac{6}{j}+1\right)\left(\frac{2}{j}+1\right)} \cdot \alpha^{4 / j} .
\end{aligned}
$$

Now $E^{[j]}\left(\bar{\zeta}_{\alpha}\right) \leq 2^{4 / j} C_{5} \alpha^{4 / j}+0+2 \cdot \frac{(2 j)^{2 / j}}{(6 / j+1)(2 / j+1)} \cdot \alpha^{4 / j}=: C_{4} \alpha^{4 / j}$.

Proposition 7 (Energy of a segment). There are constants $c_{0} \in\left(0, c_{4}\right], C_{0}<\infty$ depending only on $j>0$ such that

$$
E^{[j]}\left(\kappa_{\alpha}\right) \leq C_{0} \alpha^{4 / j} \quad \text { for any } \alpha \in\left(0, c_{0}\right] .
$$

Proof. For $\alpha \in\left(0, c_{4}\right]$ consider the curve $\bar{\zeta}_{\alpha}$ defined in Lemma 6 in arc-length parametrization. Because of

$$
\left|\kappa_{\alpha}(s)-\kappa_{\alpha}(t)\right| \geq\left|\bar{\zeta}_{2 \eta_{\alpha}}\left(s+\eta_{\alpha}\right)-\bar{\zeta}_{2 \eta_{\alpha}}\left(t+\eta_{\alpha}\right)\right|
$$

for $s, t \in\left(-\infty, \eta_{\alpha}\right]$ we obtain

$$
E^{[j]}\left(\left.\kappa_{\alpha}\right|_{\left(-\infty, \eta_{a}\right]}\right) \leq E^{[j]}\left(\bar{\zeta}_{2 \eta_{\alpha}}\right) \leq C_{4}\left(2 \eta_{\alpha}\right)^{4 / j} \stackrel{(8)}{\leq} 8^{4 / j} C_{4} \alpha^{4 / j} \quad \text { for } \alpha \leq c_{0}:=\frac{1}{8} c_{4} .
$$

By symmetry we obtain the same estimate for $E^{[j]}\left(\left.\kappa_{\alpha}\right|_{\left[-\eta_{\alpha}, \infty\right)}\right)$.

It remains to treat the case

$$
\int_{-\infty}^{-\eta_{\alpha}} \int_{\eta_{\alpha}}^{\infty}\left(\frac{1}{\left|\kappa_{\alpha}(s)-\kappa_{\alpha}(t)\right|^{j}}-\frac{1}{(s-t)^{j}}\right)^{2 / j} \mathrm{~d} s \mathrm{~d} t \quad \text { resp. } \quad \int_{\eta_{\alpha}}^{\infty} \int_{-\infty}^{-\eta_{\alpha}} \cdots .
$$

Because of $\left|\kappa_{\alpha}(s)-\kappa_{\alpha}(t)\right|=(s-t)-2\left(\eta_{\alpha}-\xi_{\alpha}\right)$ for $s,(-t) \in\left[\eta_{\alpha}, \infty\right)$, the first integral equals

$$
\begin{aligned}
& \int_{-\infty}^{-\xi_{\alpha}} \int_{\xi_{\alpha}}^{\infty}\left(\frac{1}{(s-t)^{j}}-\frac{1}{\left(s-t+2\left(\eta_{\alpha}-\xi_{\alpha}\right)\right)^{j}}\right)^{2 / j} \mathrm{~d} s \mathrm{~d} t \\
& \stackrel{(4)}{\leq} C_{2}\left(\frac{\eta_{\alpha}-\xi_{\alpha}}{\xi_{\alpha}}\right)^{2 / j} \stackrel{(7)-(9)}{\leq} 2^{2 / j} C_{2} \alpha^{4 / j} .
\end{aligned}
$$

By symmetry we arrive at the same estimate for the second integral. Summing up we conclude $E\left(\kappa_{\alpha}\right) \leq 2 \cdot 8^{4 / j} C_{4} \alpha^{4 / j}+2 \cdot 2^{2 / j} C_{2} \alpha^{4 / j}=: C_{0} \alpha^{4 / j}$.

Proposition 8 (Joining curves). Let $j \geq 2, L \geq \Lambda \geq \frac{1}{2} \Lambda \geq \ell \geq \lambda>0, \beta \in[0,2 \pi)$. Let $\gamma:[-L-\Lambda, L+\Lambda] \rightarrow \mathbb{R}^{2}$ be an open curve in arc-length parametrization satisfying

$$
\begin{array}{rlll}
\gamma(t) & =\binom{\cos \beta}{\sin \beta} t & \text { for all }|t| \leq \Lambda \quad \text { and } \\
|\gamma(t)| & \Lambda & \text { for all }|t| \geq \Lambda .
\end{array}
$$

Let $\delta:[-(\Lambda-\lambda)-\ell,(\Lambda-\lambda)+\ell] \rightarrow \mathbb{R}^{2}$ be an open curve in arc-length parametrization satisfying

$$
\begin{array}{rll}
\delta(t) & =\binom{\cos \beta}{\sin \beta}(t+(\ell-\lambda)) & \text { for all } t \leq-\ell, \\
|\delta(t)| & \leq & \text { for all }|t| \leq \ell, \quad \text { and } \\
\delta(t) & =\binom{\cos \beta}{\sin \beta}(t-(\ell-\lambda)) & \text { for all } t \geq \ell .
\end{array}
$$

Then the energy of the open arc-length parametrized curve

$$
\begin{aligned}
\gamma \curvearrowleft \delta: & {[-L-(\Lambda-\lambda)-\ell, L+(\Lambda-\lambda)+\ell] \longrightarrow \mathbb{R}^{2} } \\
t & \longmapsto \begin{cases}\gamma(t+\ell-\lambda) & \text { for } t \in[-L-(\Lambda-\lambda)-\ell,-\ell], \\
\delta(t) & \text { for } t \in[-\ell, \ell], \\
\gamma(t-\ell+\lambda) & \text { for } t \in[\ell, L+(\Lambda-\lambda)+\ell],\end{cases}
\end{aligned}
$$

fulfills the estimate $E^{[j]}(\gamma \curvearrowleft \delta) \leq E^{[j]}(\gamma)+E^{[j]}(\delta)+2 C_{2}\left(\frac{\ell-\lambda}{\lambda}\right)^{2 / j}+32 \frac{L \ell}{\Lambda^{2}}$.
In the proof we will need the fact that, for $a \geq b \geq 0, p \in[0,1]$,

$$
\begin{equation*}
a^{p}-b^{p} \leq(a-b)^{p} \tag{11}
\end{equation*}
$$

holds, which can easily be verified by showing that $\sigma^{p}+(1-\sigma)^{p}$ is monotone increasing in $\sigma$ on $\left[0, \frac{1}{2}\right]$. The fact that there is no similar estimate for $p>1$ is the reason for the restriction to $j \geq 2$.

Proof. Let

$$
\begin{aligned}
& \mathcal{A}:=[-L-(\Lambda-\lambda)-\ell,-(\Lambda-\lambda)-\ell], \\
& \mathcal{B} \\
& \mathcal{C} \\
& \mathcal{D} \\
& \mathcal{D} \\
& \mathcal{E}
\end{aligned}:=[-=[(\Lambda,(\Lambda),(\Lambda-\lambda)+\ell,-\ell], .
$$

We divide the integration domain into 25 subdomains. Because of symmetry we may restrict ourselves to 15 of them (see table on the right), but we have to keep in mind to multiply by 2 off the diagonal. The term $E^{[j]}(\gamma)$ absorbs cases $1,2,6,13,14,15, E^{[j]}(\delta)$ treats cases $7,8,10,11$; in both cases the symmetric domains are already considered.
Now we consider the cases $4,5,8,9$.

|  | $\mathcal{A}$ | $\mathcal{B}$ | $\mathcal{C}$ | $\mathcal{D}$ | $\mathcal{E}$ |
| :---: | :---: | :---: | :---: | :---: | :---: |
| $\mathcal{A}$ | 1 | 2 | 3 | 4 | 5 |
| $\mathcal{B}$ |  | 6 | 7 | 8 | 9 |
| $\mathcal{C}$ |  |  | 10 | 11 | 12 |
| $\mathcal{D}$ |  |  |  | 13 | 14 |
| $\mathcal{E}$ |  |  |  |  | 15 |

$$
\begin{aligned}
& \iint_{(\mathcal{A} \cup \mathcal{B}) \times(\mathcal{D} \cup \mathcal{E})}\left(\frac{1}{|(\gamma \curvearrowleft \delta)(s)-(\gamma \curvearrowleft \delta)(t)|^{j}}-\frac{1}{(s-t)^{j}}\right)^{2 / j} \mathrm{~d} s \mathrm{~d} t \\
= & \int_{-L-\Lambda+\lambda-\ell}^{-\ell} \int_{\ell}^{L+\Lambda-\lambda+\ell}\left(\frac{1}{|\gamma(s-\ell+\lambda)-\gamma(t+\ell-\lambda)|^{j}}-\frac{1}{(s-t)^{j}}\right)^{2 / j} \mathrm{~d} s \mathrm{~d} t \\
= & \int_{-L-\Lambda}^{-\lambda} \int_{\lambda}^{L+\Lambda}\left(\frac{1}{|\gamma(s)-\gamma(t)|^{j}}-\frac{1}{(s-t+2(\ell-\lambda))^{j}}\right)^{2 / j} \mathrm{~d} s \mathrm{~d} t
\end{aligned}
$$

The $\int_{-L-\Lambda}^{-\lambda} \int_{\lambda}^{L+\Lambda}$ part of $E^{[j]}(\gamma)$ was not already used in the absorbing process, so referring to (11) it suffices to examine

$$
\begin{aligned}
\int_{-L-\Lambda}^{-\lambda} & \int_{\lambda}^{L+\Lambda}\left(\frac{1}{|\gamma(s)-\gamma(t)|^{j}}-\frac{1}{(s-t+2(\ell-\lambda))^{j}}\right)^{2 / j} \mathrm{~d} s \mathrm{~d} t \\
& -\int_{-L-\Lambda}^{-\lambda} \int_{\lambda}^{L+\Lambda}\left(\frac{1}{|\gamma(s)-\gamma(t)|^{j}}-\frac{1}{(s-t)^{j}}\right)^{2 / j} \mathrm{~d} s \mathrm{~d} t \\
\stackrel{(11)}{\leq} & \int_{-L-\Lambda}^{-\lambda} \int_{\lambda}^{L+\Lambda}\left(\frac{1}{(s-t)^{j}}-\frac{1}{(s-t+2(\ell-\lambda))^{j}}\right)^{2 / j} \mathrm{~d} s \mathrm{~d} t \stackrel{(4)}{\leq} C_{2}\left(\frac{\ell-\lambda}{\lambda}\right)^{2 / j} .
\end{aligned}
$$

For case 3 let $s \in \mathcal{A}, t \in C$. Now $|(\gamma \curvearrowleft \delta)(s)-(\gamma \curvearrowleft \delta)(t)| \geq \Lambda-\lambda$ leads to

$$
\int_{\mathcal{A}} \int_{C}\left(\frac{1}{|(\gamma \curvearrowleft \delta)(s)-(\gamma \curvearrowleft \delta)(t)|^{j}}-\frac{1}{(s-t)^{j}}\right)^{2 / j} \mathrm{~d} s \mathrm{~d} t \leq \frac{|\mathcal{A}||C|}{(\Lambda-\lambda)^{2}} \leq 8 \frac{L \ell}{\Lambda^{2}}
$$

Since we arrive at the same situation in the remaining case 12 , we obtain $2 \cdot 8 L \ell / \Lambda^{2}$ together for both cases. Summing up and remembering that we have to multiply by 2 off the diagonal, we are finally lead to the formula stated above.

Proposition 9 (Extending a curve). Let $j, L^{*}, \Lambda^{*}, \ell^{*}, \lambda^{*}, \delta^{*}$ as in Proposition 8 , and $\beta^{*}:=0$ and $\gamma^{*}(t):=\binom{1}{0}$ tfor all $t \in\left[-L^{*}-\Lambda^{*}, L^{*}+\Lambda^{*}\right]$. Then there is a constant $C_{6}>0$ depending only on $j$ such that

$$
E^{[j]}\left(\gamma^{*} \curvearrowleft \delta^{*}\right) \leq E^{[j]}\left(\delta^{*}\right)+2 C_{2}\left(\frac{\ell^{*}-\lambda^{*}}{\lambda^{*}}\right)^{2 / j}+4 C_{6}\left(\frac{2 \ell^{*}}{\Lambda^{*}-\lambda^{*}}\right)^{1+2 / j}
$$

Of course, Propositions 8 and 9 are also true for an analogous situation in $\mathbb{R}^{n}$.
Proof. We treat cases $1,2,4,5, \ldots, 11,13,14,15$ as in the proof of Proposition 8. In case 3 (and analogously in case 12) we obtain for $s \in \mathcal{A}, t \in C$ the estimates

$$
\left|\left(\gamma^{*} \curvearrowleft \delta^{*}\right)(s)-\left(\gamma^{*} \curvearrowleft \delta^{*}\right)(t)\right| \geq|s|-\ell^{*} \quad \text { and } \quad|s-t| \leq|s|+\ell^{*}
$$

This leads to

$$
\begin{aligned}
& \int_{\mathcal{A}} \int_{C}\left(\frac{1}{\left|\left(\gamma^{*} \curvearrowleft \delta^{*}\right)(s)-\left(\gamma^{*} \curvearrowleft \delta^{*}\right)(t)\right|^{j}}-\frac{1}{(s-t)^{j}}\right)^{2 / j} \mathrm{~d} t \mathrm{~d} s \\
& \quad \leq|C| \int_{-L^{*}-\left(\Lambda^{*}-\lambda^{*}\right)-\ell^{*}}^{-\left(\Lambda^{*}-\lambda^{*}\right)-\ell^{*}}\left(\frac{1}{\left(-s-\ell^{*}\right)^{j}}-\frac{1}{\left(-s+\ell^{*}\right)^{j}}\right)^{2 / j} \mathrm{~d} s \\
& \leq 2 \ell \int_{\Lambda-\lambda}^{\infty}\left(\frac{1}{s^{j}}-\frac{1}{(s+2 \ell)^{j}}\right)^{2 / j} \mathrm{~d} s=2 \ell \int_{\Lambda-\lambda}^{\infty} \frac{1}{s^{2}}\left[1-\left(\frac{s}{s+2 \ell}\right)^{j}\right]^{2 / j} \mathrm{~d} s \\
& \quad(6) \quad(2 \ell)^{1+2 / j} j^{2 / j} \int_{\Lambda-\lambda}^{\infty} s^{-2-2 / j} \mathrm{~d} s=\frac{j^{2 / j}}{1+\frac{2}{j}} \cdot\left(\frac{2 \ell}{\Lambda-\lambda}\right)^{1+2 / j}=: \quad C_{6}\left(\frac{2 \ell}{\Lambda-\lambda}\right)^{1+2 / j} .
\end{aligned}
$$

Notice that $E^{[j]}\left(\gamma^{*}\right)=0$.
Proof of Theorem 2. For $j, x, y>0$ the function $j \mapsto\left(x^{-j}-(x+y)^{-j}\right)^{2 / j}$ is monotone increasing since, for $0<j_{1}<j_{2}$,

$$
\left(\frac{1}{x^{j_{2}}}-\frac{1}{(x+y)^{j_{2}}}\right)^{\frac{2}{j_{2}}}=\left(\frac{1}{x^{j_{2}}}-\frac{1}{(x+y)^{j_{2}}}\right)^{\frac{2}{j_{1} \cdot \frac{j_{1}}{j_{2}}}} \stackrel{(11)}{\geq}\left(\frac{1}{x^{j_{2} \cdot \frac{j_{1}}{j_{2}}}}-\frac{1}{(x+y)^{j_{2} \cdot \frac{j_{1}}{j_{2}}}}\right)^{\frac{2}{j_{1}}}
$$

This yields that $E^{[j]}$ is monotone increasing in $j$, so we may restrict ourselves to $j \in$ $[2,4)$. Fix $k_{0} \in \mathbb{N}$, such that

$$
\begin{equation*}
k_{0}>\max \left(\frac{1}{c_{0}}, 1+\left(\frac{C_{7}}{\varepsilon}\right)^{c_{7}}\right) \tag{12a}
\end{equation*}
$$

for constants $c_{7}>0, C_{7}<\infty$ that depend only on $j$ and will be defined later in this proof. Let

$$
\begin{equation*}
L_{k_{0}}:=2\left(1+\frac{4}{3} \frac{1}{k_{0}-1}\right) \sqrt{k_{0}-1} \quad>\quad \Lambda_{k_{0}}:=1 \tag{12b}
\end{equation*}
$$

and define for $k \in \mathbb{N}, k \geq k_{0}$, the following positive quantities, whose estimates will be proven by induction.

$$
\begin{align*}
\alpha_{k} & :=\frac{1}{k}  \tag{12c}\\
\beta_{k} & :=\sum_{k=k_{0}}^{k-1} \frac{1}{\kappa},  \tag{12d}\\
r_{k} & :=\frac{\Lambda_{k}^{2} \alpha_{k}^{2}}{2 L_{k} \eta_{\alpha_{k}}} \leq \frac{\alpha_{k}}{6} \leq \frac{1}{6},  \tag{12e}\\
\ell_{k} & :=r_{k} \eta_{k_{0}}:=0,  \tag{12f}\\
\eta_{k} & \frac{1}{2} \Lambda_{k},
\end{align*}
$$

$$
\begin{align*}
\lambda_{k} & :=r_{k} \xi_{\alpha_{k}} \leq \ell_{k},  \tag{12~g}\\
L_{k+1} & :=L_{k}+\left(\Lambda_{k}-\lambda_{k}\right)+\ell_{k}-\Lambda_{k+1} \geq L_{k} \geq 1,  \tag{12h}\\
\Lambda_{k+1} & :=r_{k} y\left(\alpha_{k}\right) \leq \ell_{k} \leq 1 . \tag{12i}
\end{align*}
$$

The estimate in (12g), the first one in (12i) and the last one in (12e) are immediate.
For the initial step $k=k_{0}$ the inequality $r_{k_{0}}<\alpha_{k_{0}}^{2} /\left(2 \cdot 3 \alpha_{k_{0}}\right)$ verifies (12e),

$$
\ell_{k_{0}}=\frac{\Lambda_{k_{0}}^{2} \alpha_{k_{0}}^{2}}{2 L_{k_{0}}} \leq \frac{\Lambda_{k_{0}}}{2}=\frac{1}{2}<1
$$

implies (12f) and the last inequality in (12i). $L_{k_{0}+1}-L_{k_{0}}=\left(\Lambda_{k_{0}}-\Lambda_{k_{0}+1}\right)+\left(\ell_{k_{0}}-\lambda_{k_{0}}\right) \geq$ $0+0$ yields (12h).

For the step $k \rightarrow k+1$ we obtain as above

$$
\begin{aligned}
r_{k+1} & =\frac{\Lambda_{k+1}^{2} \alpha_{k+1}^{2}}{2 L_{k+1} \eta_{\alpha_{k+1}}^{2}} \stackrel{(8)}{\leq} \frac{\alpha_{k+1}^{2}}{6 \alpha_{k+1}} \leq \frac{\alpha_{k+1}}{6}, \\
\ell_{k+1} & =\frac{\Lambda_{k+1}^{2} \alpha_{k+1}^{2}}{2 L_{k+1}} \leq \frac{\Lambda_{k+1}}{2}, \\
L_{k+2}-L_{k+1} & =\underbrace{\Lambda_{k+1}-\Lambda_{k+2}}_{\geq \Lambda_{k+1}-\ell_{k+1}}+\underbrace{\ell_{k+1}-\lambda_{k+1}}_{\geq 0} \geq \Lambda_{k+1}-\frac{1}{2} \Lambda_{k+1} \geq 0 .
\end{aligned}
$$

Notice that $\left(L_{k}\right)_{k \in \mathbb{N}}$ is a monotone increasing sequence; we will see in (14) that it converges, whereas $\left(\beta_{k}\right)_{k \in \mathbb{N}}$ is the diverging harmonic sequence. All other sequences are monotone decreasing sequences that converge to 0 .
We consider the sequence of functions

$$
\begin{gathered}
\delta_{k}:\left[-\Lambda_{k}+\lambda_{k}-\ell_{k}, \Lambda_{k}-\lambda_{k}+\ell_{k}\right] \rightarrow \mathbb{R}^{2}, \\
t \mapsto\left(\begin{array}{c}
\cos \beta_{k}-\sin \beta_{k} \\
\sin \beta_{k} \\
\cos \beta_{k}
\end{array}\right) r_{k} K_{\alpha_{k}}\left(\frac{1}{r_{k}} t\right)
\end{gathered}
$$

and start with $\gamma_{k_{0}}:[-2,2] \rightarrow \mathbb{R}^{2}, t \mapsto\binom{t}{0}$.


If we define $\gamma_{k+1}:=\gamma_{k} \curvearrowleft \delta_{k}$ for all $k \in \mathbb{N}, k \geq k_{0}$, the conditions of Proposition 2 are true in any step, since $\Lambda_{k+1} \leq y\left(\alpha_{k}\right)$ by (12e). Now, for $K \geq k_{0}$,

$$
E^{[j]}\left(\gamma_{K}\right) \leq E^{[j]}\left(\gamma_{k_{0}}\right)+\sum_{k=k_{0}}^{K} E^{[j]}\left(\delta_{k}\right)+2 C_{2} \sum_{k=k_{0}}^{K}\left(\frac{\ell_{k}-\lambda_{k}}{\lambda_{k}}\right)^{2 / j}+32 \sum_{k=k_{0}}^{K} \frac{L_{k} \ell_{k}}{\Lambda_{k}^{2}},
$$

where $E^{[j]}\left(\gamma_{k_{0}}\right)=0, E^{[j]}\left(\delta_{k}\right) \leq C_{0} / k^{4 / j}$ by Proposition 7,

$$
\begin{equation*}
\frac{\ell_{k}-\lambda_{k}}{\lambda_{k}} \stackrel{(7)}{\leq} \frac{\eta_{\alpha_{k}}-\xi_{\alpha_{k}}}{3 \cdot \frac{1}{2} \alpha_{k}} \stackrel{(9)}{<} \frac{2}{k^{2}}, \tag{13}
\end{equation*}
$$

$L_{k} \ell_{k} / \Lambda_{k}^{2}=1 /\left(2 k^{2}\right)$ by (12e), and therefore

$$
E^{[j]}\left(\gamma_{K}\right)<\left(C_{0}+2^{2 / j+1} C_{2}+16\right) \sum_{k=k_{0}}^{K} \frac{1}{k^{4 / j}} \leq \underbrace{\frac{C_{0}+2^{2 / j+1} C_{2}+16}{\frac{4}{j}-1}}_{=: C_{8}} \cdot \frac{1}{\left(k_{0}-1\right)^{4 / j-1}}
$$

Because of

$$
\begin{align*}
& L_{K}-L_{k_{0}} \leq \sum_{k=k_{0}}^{K} \ell_{k}+\sum_{k=k_{0}}^{K} \Lambda_{k} \stackrel{(12 \mathrm{i})}{\leq} 1+2 \sum_{k=k_{0}}^{K+1} \ell_{k}  \tag{14}\\
&(8),(12 e),(12 \mathrm{f}) \\
& \leq 1+2 \sum_{k=k_{0}}^{K+1} \frac{\alpha_{k}}{6} \cdot 4 \alpha_{k} \leq 1+\frac{4}{3} \sum_{k=k_{0}}^{K+1} \frac{1}{k^{2}} \leq 1+\frac{4}{3} \frac{1}{k_{0}-1}
\end{align*}
$$

the limit $L_{\infty}:=\lim _{K \rightarrow \infty} L_{K}<\infty$ exists.
The reparametrizations $\tilde{\gamma}_{k}:[-1,1] \rightarrow \mathbb{R}^{2}, \tilde{\gamma}_{k}(t):=\gamma_{k}\left(\left(L_{k}+\Lambda_{k}\right) t\right)$, of $\gamma_{k}$ to constant velocity (depending on $k$ ) form a $C^{0}$-Сauchy sequence by (14) and $\Lambda_{k} \searrow 0$ that converges to some limit curve $\tilde{\gamma}_{\infty} \in C^{0}\left([-1,1], \mathbb{R}^{2}\right)$. Now, by Lemma 15 , we obtain $E^{[j]}\left(\tilde{\gamma}_{\infty}\right) \leq \liminf _{k \rightarrow \infty} E^{[j]}\left(\gamma_{k}\right)<\infty$. Let $\gamma_{\infty}:\left[-L_{\infty}, L_{\infty}\right] \rightarrow \mathbb{R}^{2}, t \mapsto \tilde{\gamma}\left(t / L_{\infty}\right)$.
The proof that $\gamma_{\infty}$ is not differentiable at 0 is deferred to Proposition 11.
Now we are going to extend $\gamma_{\infty}:\left[-L_{\infty}, L_{\infty}\right] \rightarrow \mathbb{R}^{2}$ to $\mathbb{R} \rightarrow \mathbb{R}^{2}$ via Proposition 9. Let $\Lambda^{*}:=L_{k_{0}}+\Lambda_{k_{0}}=L_{k_{0}}+1, \ell^{*}:=L_{\infty}-L_{k_{0}}, \lambda^{*}:=\Lambda_{k_{0}}=1, \delta^{*}:=\gamma_{\infty}$, and $L^{*} \geq \Lambda^{*}$ arbitrary. The fact that our estimates will not depend on $L^{*}$ will allow us to finally take $L^{*} \rightarrow \infty$.
We compute $\ell^{*}=L_{\infty}-L_{k_{0}} \stackrel{(14)}{\leq} 1+\frac{4}{3} \frac{1}{k_{0}-1}$ and

$$
\frac{2 \ell^{*}}{\Lambda^{*}-\lambda^{*}} \stackrel{(12 \mathrm{~b})}{=} \frac{2\left(1+\frac{4}{3} \frac{1}{k_{0}-1}\right)}{2\left(1+\frac{4}{3} \frac{1}{k_{0}-1}\right) \sqrt{k_{0}-1}}=\frac{1}{\sqrt{k_{0}-1}} .
$$

Summing up and choosing $c_{7}:=1 / \min \left(\frac{4}{j}-1, \frac{2}{j}, \frac{1}{2}+\frac{1}{j}\right)<\infty$, we obtain

$$
\begin{aligned}
E^{[j]}\left(\gamma^{*} \curvearrowleft \delta^{*}\right) & \leq E^{[j]}\left(\delta^{*}\right)+2 C_{2}\left(\frac{\ell^{*}-\lambda^{*}}{\lambda^{*}}\right)^{2 / j}+4 C_{6}\left(\frac{2 \ell^{*}}{\Lambda^{*}-\lambda^{*}}\right)^{1+2 / j} \\
& \leq \frac{C_{8}}{\left(k_{0}-1\right)^{4 / j-1}}+2 C_{2}\left(\frac{\frac{4}{3}}{k_{0}-1}\right)^{2 / j}+\frac{4 C_{6}}{\left(k_{0}-1\right)^{1 / 2+1 / j}} \\
& \leq\left(C_{8}+2\left(\frac{4}{3}\right)^{2 / j} C_{2}+4 C_{6}\right) \cdot \frac{1}{\left(k_{0}-1\right)^{1 / c_{7}}} \\
& =: \frac{C_{7}}{\left(k_{0}-1\right)^{1 / c 7}} \stackrel{(12 a)}{<} \varepsilon .
\end{aligned}
$$

## Remark 10.

(i) In the previous proof the limit curve $\gamma_{\infty}$ was defined as

$$
\gamma_{\infty}(t):=\lim _{k \rightarrow \infty} \gamma_{k}\left(\frac{L_{k}-\Lambda_{k}}{L_{\infty}} t\right) .
$$

The fact that all curves $\gamma_{k}$ are parametrized by arc-length yields for $t>0$

$$
\begin{equation*}
\gamma_{\infty}(t)=\lim _{\substack{k \rightarrow \infty \\ L_{k}+\Lambda_{k}>t}} \gamma_{k}(t)=\gamma_{K}\left(t-\sum_{k=K}^{\infty}\left(\ell_{k}-\lambda_{k}\right)\right) \tag{15}
\end{equation*}
$$

where $K$ is chosen so large that $\Lambda_{K}<t$. The second identity is due to the fact that

$$
\gamma_{k}(t)=\gamma_{k+1}\left(t+\left(\ell_{k}-\lambda_{k}\right)\right)
$$

for all $k \geq k_{0}$ and $\lambda_{k} \leq t \leq L_{k}+\Lambda_{k}$.
(ii) Because of $r_{k+1} \leq \frac{1}{6} \Lambda_{k+1}^{2} \alpha_{k} \leq(6 k)^{-2} r_{k}$ the sequence $\left(r_{k}\right)_{k \in \mathbb{N}}$ decreases "superexponentially".

Proposition 11. The curve $\gamma_{\infty}$ constructed in the proof of Theorem 2 is not differentiable at 0.

Proof. Let $\tau_{k}:=\ell_{k}+\sum_{k=k}^{\infty}\left(\ell_{\kappa}-\lambda_{k}\right) \in\left(0, L_{\infty}\right)$. Since $\tau_{k} \geq \ell_{k} \geq \lambda_{k}$ we infer from (15) that $\gamma_{\infty}\left(\tau_{k}\right)=\gamma_{k}\left(\ell_{k}\right)$ and note $\Varangle\left(\gamma_{\infty}\left(\tau_{k}\right), e_{1}\right) \equiv \beta_{k} \bmod 2 \pi$, where $\Varangle(\cdot, \cdot) \in[0,2 \pi)$, $e_{1}:=\binom{1}{0}$ and $x \equiv y \bmod 2 \pi$ iff there is an $m \in \mathbb{Z}$ with $y-x=2 \pi m$. We compute

$$
\begin{aligned}
& \tau_{k} \leq 2 \sum_{k=k}^{\infty} \ell_{\kappa} \\
& \stackrel{(14)}{\leq} \frac{4}{3} \frac{1}{k-1} \searrow 0 \\
& \xi_{\alpha_{k}} \stackrel{(7)}{\geq} \\
& \frac{3}{2} \alpha_{k} \stackrel{(9)}{\Longrightarrow} \\
& \frac{\ell_{k}}{\lambda_{k}}=\frac{\eta_{\alpha_{k}}}{\xi_{\alpha_{k}}} \stackrel{(8)}{\leq} \frac{\xi_{\alpha_{k}}+3 \alpha_{k}^{3}}{\xi_{\alpha_{k}}} \leq 1+\frac{2}{k^{2}} \searrow 0
\end{aligned}
$$

and

$$
\frac{\tau_{k}}{\lambda_{k}} \leq \frac{\ell_{k}}{\lambda_{k}}+\sum_{k=k}^{\infty} \frac{\ell_{\kappa}-\lambda_{k}}{\lambda_{\kappa}} \stackrel{(13)}{<} \frac{\ell_{k}}{\lambda_{k}}+2 \sum_{\kappa=k}^{\infty} \frac{1}{\kappa^{2}} \leq 1+\frac{2}{k^{2}}+\frac{2}{k-1} \searrow 1
$$

for $k \rightarrow \infty$. Since $\left|\gamma_{k}\left(\ell_{k}\right)\right|=\lambda_{k}$ and $\gamma_{\infty}(0)=0$, we arrive at

$$
\cos \beta_{k}=\cos \Varangle\left(\gamma_{\infty}\left(\tau_{k}\right), e_{1}\right)=\frac{\left\langle\gamma_{\infty}\left(\tau_{k}\right), e_{1}\right\rangle}{\left|\gamma_{\infty}\left(\tau_{k}\right)\right|}=\frac{\tau_{k}}{\lambda_{k}}\left\langle\frac{\gamma_{\infty}\left(\tau_{k}\right)-\gamma_{\infty}(0)}{\tau_{k}}, e_{1}\right\rangle
$$

If $\gamma_{\infty}$ were differentiable at 0 , the right hand side would converge. But the left hand side does not converge since the sequence $\left(\beta_{k} \bmod 2 \pi\right)_{k \in \mathbb{N}, k \geq k_{0}}$ is dense in $[0,2 \pi)$, as we will see now.
For any $\delta>0$ choose $K>\max \left(k_{0}, \frac{1}{\sqrt{\delta}}\right)$. Since the harmonic series does not converge and $\frac{1}{k}-\frac{1}{k+1} \in(0, \delta)$ for all $k \geq K$, for any $y \in\left[\beta_{K}, \infty\right)$ there is some $k_{y} \geq K$ with $\beta_{k_{y}}-y \in[0, \delta)$. So, for any $x \in[\delta, 2 \pi-\delta]$ we may find some $y \in\left[\beta_{K}, \infty\right)$ with $x \equiv y \bmod 2 \pi$, which implies $\left|\left(\beta_{k_{y}}-2 \pi m\right)-x\right|=\beta_{k_{y}}-y<\delta$ for some $m \in \mathbb{Z}$.

We will proceed to the situation of closed curves.
For the Möbius energy $(j, p)=(2,1)$ we obtain our result immediately by applying the Möbius invariance. Let $\varepsilon>0$. Carrying out an inversion on a circle whose center does not lie on the curve referred to by Theorem 2 we obtain a closed curve of energy $4+\varepsilon$ according to [FHW94, Thm. 2.1]. Remember that the Möbius energy of a circle amounts to 4 .
If we just want to obtain a closed curve of finite energy that is not differentiable at one point, i. e. skipping the "infimal" property, we might extend the curve constructed in Theorem 2 as indicated in the drawing on the right, using estimates as in Proposition 8.


The basic idea in proving the "infimal" property is to "glue" the curve constructed in the foregoing proofs on a cylinder; for a sphere we would arrive at the same situation. As we mentioned in the introduction our technique applies to a wide range of open curves, more precisely to planar curves $\delta$ that are admittable in sense of Proposition 8.

Proposition 12 (Projecting a curve onto a cylinder). For $j \in(0,4), \ell \in\left(0, \frac{1}{4}\right]$ let $\gamma:[-\ell, \ell] \rightarrow \mathbb{R}^{2}$ be a curve in arc-length parametrization, $E^{[j]}(\gamma)<\infty$. Then the energy of the projection $\gamma^{*}:[-\ell, \ell] \rightarrow \mathbb{R}^{3}$ of $\gamma$ onto a cylinder of radius $\frac{1}{2 \pi}$ satisfies

$$
E^{[j]}(\gamma) \leq E^{[j]}\left(\gamma^{*}\right) \leq E^{[j]}(\gamma)+\omega\left(\ell, E^{[j]}(\gamma)\right),
$$

where $\omega:\left[0, \frac{1}{4}\right] \times[0, \infty) \rightarrow[0, \infty)$ is a continuous function with $\omega(0, \cdot) \equiv 0$ depending only on $j$.

By projection we mean that the plane in which $\gamma$ lies is "glued" onto the cylinder, so lengths are preserved, i. e. $D_{\gamma^{*}}=D_{\gamma}$, and $\operatorname{since} \sin \pi x \geq \pi x-2 \pi x^{3}$ on $(0, \infty)$ we obtain

$$
\begin{equation*}
\left|\gamma^{*}(s)-\gamma^{*}(t)\right| \geq \frac{1}{\pi} \sin (\pi|\gamma(s)-\gamma(t)|) \geq|\gamma(s)-\gamma(t)|-2|\gamma(s)-\gamma(t)|^{3} \tag{16}
\end{equation*}
$$

This result can be extended to $j p \geq 2$, but due to the lack of the bi-Lipschitz property (19) our proof fails for $j p<2$.
In the proof of Theorem 2 we used the fact that the energy of a line amounts to zero for all $j$ together with the monotonicity of $j \mapsto E^{[j]}$ to transfer the "infimal" property from
$j \in[2,4)$ to $(0,4)$. In the case of closed curves we face the problem that the $\left(j, \frac{2}{j}\right)$ energy of the circle depends on $j$. So, for $j \in(0,2)$ we will make use of Minkowski's inequality. Applying the mean value theorem to $x \mapsto x^{1 / p}$, we arrive at

$$
\begin{equation*}
a-b \leq a^{1-1 / p} p\left(a^{1 / p}-b^{1 / p}\right) \text { for } \quad a \geq b \geq 1, \quad p \in[1, \infty) . \tag{17}
\end{equation*}
$$

Proof. The first inequality is an immediate consequence of (16). For the second we start computing for $x \in\left(0, \frac{1}{2}\right]$

$$
\begin{align*}
\frac{1}{\left(x-2 x^{3}\right)^{j}}-\frac{1}{x^{j}} & =\frac{1}{\left(x-2 x^{3}\right)^{j}}\left(1-\left(1-2 x^{2}\right)^{j}\right) \stackrel{(5)}{\leq} \frac{1}{x^{j}\left(1-2 x^{2}\right)^{j}} \cdot \frac{2 j x^{2}}{1-2 x^{2}} \\
& =2 j \cdot \frac{x^{2-j}}{\left(1-2 x^{2}\right)^{1+j}} \leq 2^{2+j} j x^{2-j} . \tag{18}
\end{align*}
$$

If $j \leq 2$ we obtain using the Minkowski inequality

$$
\begin{aligned}
E^{[j]}\left(\gamma^{*}\right)^{j / 2}-E^{[j]}(\gamma)^{j / 2} & \leq\left[\iint_{[-\ell, \ell]^{2}}\left(\frac{1}{\left|\gamma^{*}(s)-\gamma^{*}(t)\right|^{j}}-\frac{1}{|\gamma(s)-\gamma(t)|^{j}}\right)^{2 / j} \mathrm{~d} s \mathrm{~d} t\right]^{j / 2} \\
& \leq(2 \ell),(18) \\
\leq & \left.(2 \ell)^{2}\left(2^{2+j} j(2 \ell)^{2-j}\right)^{2 / j}\right]^{j / 2} \leq 2^{4+j} j \ell^{2},
\end{aligned}
$$

which implies $E^{[j]}\left(\gamma^{*}\right)-E^{[j]}(\gamma) \leq\left(E^{[j]}(\gamma)^{j / 2}+2^{j} j\right)^{2 / j-1} \cdot 2^{5+j} \ell^{2}$ by (17).
In case $j \in(2,4)$ we arrive at

$$
\left.\begin{array}{rl}
E^{[j]}\left(\gamma^{*}\right)-E^{[j]}(\gamma) & \stackrel{(11)}{\leq} \\
& \stackrel{(16),(18)}{\leq}
\end{array} \int_{[-\ell, \ell]^{2}}\left(\frac{1}{\left|\gamma^{*}(s)-\gamma^{*}(t)\right|^{j}}-\frac{1}{|\gamma(s)-\gamma(t)|^{j}}\right)^{2 / j} \iint_{[-\ell, \ell]^{2}}|\gamma(s)-\gamma(t)|^{4 / j-2} \mathrm{~d} s \mathrm{~d} t \mathrm{~d} t\right)
$$

Proposition 13 (Modifying a closed curve). Let $\zeta: \mathbb{R} / \mathbb{Z} \rightarrow \mathbb{R}^{3}$ be a circle of length 1 parametrized in arc-length by $t \mapsto \frac{1}{2 \pi}(\sin (2 \pi t), 1-\cos (2 \pi t), 0)$, and $\lambda \in\left(0, \frac{1}{8}\right], \ell>0$, $j \in(0,4)$. If we change $\zeta$ on $(-\lambda, \lambda)$ by inserting a curve segment (parametrized by constant velocity) of length $2 \ell$ that is completely contained in

$$
\left\{\left.\frac{1}{2 \pi}(\sin (2 \pi t), 1-\cos (2 \pi t), u) \right\rvert\, t \in[-\lambda, \lambda], u \in \mathbb{R}\right\} \cap B_{\sin (\pi \lambda) / \pi}(0),
$$

which is a ball in the cylinder induced by $\zeta$, the new curve $\tilde{\zeta}: \mathbb{R} / \mathbb{Z} \rightarrow \mathbb{R}^{3}$ satisfies

$$
\left|E^{[j]}(\tilde{\zeta})-E^{[j]}(\zeta)\right| \leq E^{[j]}\left(\left.\tilde{\zeta}\right|_{\left[-\frac{1}{4}, \frac{1}{4}\right]}\right)+\tilde{\omega}\left(\ell+\left(\frac{\ell}{\lambda}-1\right)\right),
$$

where $\tilde{\omega}:[0, \infty) \rightarrow[0, \infty)$ is a continuous function with $\tilde{\omega}(0)=0$ depending only on $j$.

For $j \leq 2$ the estimate $E^{[j]}(\zeta) \leq E^{[j]}(\tilde{\zeta})$ immediately follows from [ACF ${ }^{+} 03$, Cor. 3] implying $E^{[j]}(\zeta) \leq E^{[j]}(\tilde{\zeta}) \leq E^{[j]}(\zeta)+E^{[j]}\left(\left.\tilde{\zeta}\right|_{\left[-\frac{1}{4}, \frac{1}{4}\right]}\right)+\tilde{\omega}\left(\ell+\left(\frac{\ell}{\lambda}-1\right)\right)$.

Proof. We will again split up the integration domain, but since we intend to treat closed curves we have to consider a parallelogram

$$
\mathcal{R}:=\left\{(s, t) \in \mathbb{R}^{2} \left\lvert\, s \in\left[-\frac{1}{2}, \frac{1}{2}\right]\right., t \in\left[s+\frac{1}{2}, s-\frac{1}{2}\right]\right\},
$$

which yields $D_{\zeta}(s, t)=|s-t|$ for $(s, t) \in \mathcal{R}$. Due to symmetry we may restrict to $s \geq 0$ and cover $\mathcal{R} \cap$ $([0, \infty) \times \mathbb{R})$ by $\mathcal{A} \cup \mathcal{B} \cup \mathcal{C} \cup \mathcal{D} \cup \mathcal{E} \cup \mathcal{F} \cup\left[-\frac{1}{4}, \frac{1}{4}\right]^{2}$, where

$$
\begin{array}{rlll}
\mathcal{A} & :=\left(\left[\left[\frac{1}{4}, \frac{1}{2}\right]\right.\right. & \times[-\lambda, \lambda]) & \cap \mathcal{R}, \\
\mathcal{B} & :=([0, \lambda] & \left.\times\left[-\frac{1}{2},-\frac{1}{4}\right]\right) & \cap \mathcal{R}, \\
\mathcal{C} & :=\left([0, \lambda] \times\left[\frac{1}{4}, \frac{1}{2}+\lambda\right]\right) & \cap \mathcal{R}, \\
\mathcal{D} & :=\left(\left[\frac{1}{2}-\lambda, \frac{1}{2}\right] \times[1-\lambda, 1]\right) & \cap \mathcal{R}, \\
\mathcal{E} & :=\left(\left[\lambda, \frac{1}{2}\right]\right. & \times[\lambda, 1-\lambda]) & \cap \mathcal{R}, \\
\mathcal{F} & :=\left(\left[\lambda, \frac{1}{2}-\lambda\right] \times\left[-\frac{1}{2}+\lambda,-\lambda\right]\right) & \cap \mathcal{R} .
\end{array}
$$

Note that $\left.\tilde{\zeta}\right|_{Y}=\left.\zeta\right|_{Y}$, where $Y=\left[-\frac{1}{2},-\lambda\right]$ or $\left[\lambda, \frac{1}{2}\right]$. Furthermore $|\dot{\tilde{\zeta}}(s)|=\frac{\ell}{\lambda}$ for $s \in[-\lambda, \lambda]$ and $|\dot{\tilde{\zeta}}(s)|=$ $|\dot{\zeta}(s)|=1$ elsewhere.


Since the subdomains $\mathcal{Y}=\mathcal{A}, \mathcal{B}, \mathcal{C}, \mathcal{D}$ are bounded away from the diagonal we may
estimate the corresponding integrals to the measure of their domain obtaining

$$
\iint_{y}\left(\frac{1}{|\tilde{\zeta}(s)-\tilde{\zeta}(t)|^{j}}-\frac{1}{D_{\tilde{\zeta}}(s, t)^{j}}\right)^{2 / j}|\dot{\tilde{\zeta}}(s)||\dot{\tilde{\zeta}}(t)| \mathrm{d} t \mathrm{~d} s \quad \leq \quad C_{9} \ell
$$

for some generic $C_{9}<\infty$. The identity $D_{\tilde{\zeta}}=D_{\zeta}$ on $\mathcal{E}$ yields

$$
\iint_{\mathcal{E}}\left(\frac{1}{|\tilde{\zeta}(s)-\tilde{\zeta}(t)|^{j}}-\frac{1}{\left.D_{\tilde{\zeta}(s, t)}\right)^{j}}\right)^{2 / j}|\dot{\tilde{\zeta}}(s)||\dot{\tilde{\zeta}}(t)| \mathrm{d} t \mathrm{~d} s=\iint_{\mathcal{E}}\left(\frac{1}{|\zeta(s)-\zeta(t)|^{j}}-\frac{1}{\left.D_{\zeta}(s, t)\right)^{j}}\right)^{2 / j} \mathrm{~d} t \mathrm{~d} s .
$$

In the remaining case $\mathcal{F}$, where $D_{\tilde{\zeta}}(s, t) \leq(s-t)+2(\ell-\lambda)$ for $(s, t) \in \mathcal{F}$, we first consider $j \leq 2$ and apply Minkowski's inequality

$$
\begin{aligned}
& {\left[\iint_{\mathcal{F}}\left(\frac{1}{|\tilde{\zeta}(s)-\tilde{\zeta}(t)|^{j}}-\frac{1}{D_{\tilde{\xi}}(s, t)^{j}}\right)^{2 / j}|\dot{\tilde{\zeta}}(s)||\dot{\tilde{\zeta}}(t)| \mathrm{d} t \mathrm{~d} s\right]^{j / 2}-\left[\iint_{\mathcal{F}}\left(\frac{1}{|\zeta(s)-\zeta(t)|^{j}}-\frac{1}{D_{\zeta}(s, t)^{j}}\right)^{2 / j} \mathrm{~d} t \mathrm{~d} s\right]^{j / 2}} \\
& \quad \leq\left[\iint_{\mathcal{F}}\left(\frac{1}{(s-t)^{j}}-\frac{1}{(s-t+2(\ell-\lambda))^{j}}\right)^{2 / j} \mathrm{~d} t \mathrm{~d} s\right]_{\stackrel{(4)}{j / 2} \quad C_{2}^{j / 2}\left(\frac{\ell}{\lambda}-1\right),}^{\leq}
\end{aligned}
$$

which implies

$$
\begin{aligned}
& \iint_{\mathcal{F}}\left(\frac{1}{|\tilde{\zeta}(s)-\tilde{\zeta}(t)|^{j}}-\frac{1}{D_{\tilde{\zeta}}(s, t)^{j}}\right)^{2 / j}|\dot{\tilde{\zeta}}(s)||\dot{\tilde{\zeta}}(t)| \mathrm{d} t \mathrm{~d} s-\iint_{\mathcal{F}}\left(\frac{1}{|\zeta(s)-\zeta(t)|^{j}}-\frac{1}{D_{\zeta}(s, t)^{j}}\right)^{2 / j} \mathrm{~d} t \mathrm{~d} s \\
& \quad \leq\left[E^{[j]}(\zeta)^{j / 2}+C_{2}^{j / 2}\left(\frac{\ell}{\lambda}-1\right)\right]^{2 / j-1} \cdot \frac{2}{j} \cdot C_{2}^{j / 2}\left(\frac{\ell}{\lambda}-1\right)
\end{aligned}
$$

by (17). For $j \in(2,4)$ we obtain

$$
\begin{aligned}
& \iint_{\mathcal{F}}\left(\frac{1}{|\tilde{\zeta}(s)-\tilde{\zeta}(t)|^{j}}-\frac{1}{D_{\tilde{\zeta}(s, t)^{j}}^{j}}\right)^{2 / j}|\dot{\tilde{\zeta}}(s)||\dot{\tilde{\zeta}}(t)| \mathrm{d} t \mathrm{~d} s-\iint_{\mathcal{F}}\left(\frac{1}{|\zeta(s)-\zeta(t)|^{j}}-\frac{1}{D_{\zeta}(s, t)^{j}}\right)^{2 / j} \mathrm{~d} t \mathrm{~d} s \\
& \quad \stackrel{(11)}{\leq} \quad \iint_{\mathcal{F}}\left(\frac{1}{(s-t)^{j}}-\frac{1}{(s-t+2(\ell-\lambda))^{j}}\right)^{2 / j} \mathrm{~d} t \mathrm{~d} s \stackrel{(4)}{\leq} \quad C_{2}\left(\frac{\ell}{\lambda}-1\right)^{2 / j},
\end{aligned}
$$

which concludes the proof of $E^{[j]}(\tilde{\zeta})-E^{[j]}(\zeta) \leq \cdots$.
For $E^{[j]}(\zeta)-E^{[j]}(\tilde{\zeta}) \leq \cdots$ we may use the same estimates as above.
Proof of Corollary 3. Let $\varepsilon>0$ and take the open curve $\gamma: \mathbb{R} \rightarrow \mathbb{R}^{2}$ that was constructed in the proof of Theorem 2 choosing $E^{[j]}(\gamma)<\varepsilon / 3$. Recall that $E^{j, p}$ is scaling invariant, so scaling down $\gamma$ we may assume that $\left.\gamma\right|_{\mathbb{R} \backslash[-\mu, \mu]}$ lies on the the $x$-axis for some $\mu \in(0, \lambda]$ without affecting its energy. We denote the length of the curve $\gamma_{[-\lambda, \lambda]}$ by $\ell$ which tends to 0 as $\lambda \searrow 0$. By choosing $\mu$ sufficiently small the ratio $\frac{\ell}{\lambda}$ of curvelength $2 \ell$ and diameter $2 \lambda$ tends to 1 . So the term $\ell+\left(\frac{\ell}{\lambda}-1\right)$ may be chosen arbitrary small, i. e the quantities $\omega\left(\ell, E^{[j]}(\gamma)\right)$ and $\tilde{\omega}\left(\ell+\left(\frac{\ell}{\lambda}-1\right)\right)$ are both bounded by $\varepsilon / 3$. Applying Propositions 12 and 13 we obtain $\left|E^{[j]}(\tilde{\zeta})-E^{[j]}(\zeta)\right|<\varepsilon$.

Lemma 14 (Bi-Lipschitz continuity of finite energy curves for $\boldsymbol{j} \boldsymbol{p} \geq \mathbf{2}$ ). Let $j p \geq 2$. There is a continuous monotone decreasing function $K=K_{j, p}:[0, \infty) \rightarrow(0, \infty)$ such that, for any finite-energy $\gamma \in C^{0,1}\left(X, \mathbb{R}^{3}\right)$ and $s, t \in X$

$$
\begin{equation*}
|\gamma(s)-\gamma(t)| \geq K\left(E^{j, p}(\gamma)\right) \cdot D_{\gamma}(s, t) . \tag{19}
\end{equation*}
$$

Moreover the energy of non-injective curves is infinite.

The proof in [O'H92b, Thm. 2.3] or [O'H03, Thm. 2.4.1(2)], which restricts to closed curves of length 1 , also holds for open curves. Since $E^{j, p}$ is invariant under scaling and reparametrization by definition, we obtain the claim for arbitrary curves of finite energy.

A curve $\gamma: X \rightarrow \mathbb{R}^{3}$ is called

- locally rectifiable iff it is rectifiable on any compact subset of $X$;
- regular absolutely continuous iff it is absolutely continuous and $\dot{\gamma} \neq 0$ almost everywhere.

Lemma 15 (Lower semi-continuity of the length and energy functional in the topology of pointwise convergence). Let $X=[a, b], a<b \in[-\infty, \infty]$ or $X=\mathbb{R} /(\ell \mathbb{Z})$, $\ell>0$. Let $\left(\gamma_{n}\right)_{n \in \mathbb{N}}$ be a sequence of locally rectifiable curves $\gamma_{n}: X \rightarrow \mathbb{R}^{3}$ that converge pointwise to a curve $\gamma_{0}: X \rightarrow \mathbb{R}^{3}$.
(i) For any connected subdomain $I \subset X$ we obtain

$$
\begin{equation*}
\mathscr{L}\left(\left.\gamma_{0}\right|_{I}\right) \leq \liminf _{n \rightarrow \infty} \mathscr{L}\left(\left.\gamma_{n}\right|_{I}\right), \tag{20}
\end{equation*}
$$

where $\mathscr{L}$ denotes the length functional. In particular $\gamma_{0}$ is also locally rectifiable if $\lim \inf _{n \rightarrow \infty} \mathscr{L}\left(\left.\gamma_{n}\right|_{I}\right)$ is finite for all compact subsets $I \subset X$.
(ii) If $\gamma_{0}, \gamma_{1}, \gamma_{2}, \ldots$ are in addition regular absolutely continuous curves parametrized by constant velocities $v_{0}, v_{1}, v_{2}, \ldots$, which are uniformly bounded, and $j, p>0$, we arrive at

$$
\begin{equation*}
E^{j, p}\left(\gamma_{0}\right) \leq \liminf _{n \rightarrow \infty} E^{j, p}\left(\gamma_{n}\right) \tag{21}
\end{equation*}
$$

(iii) If $\gamma_{0}, \gamma_{1}, \gamma_{2}, \ldots$ are merely regular absolutely continuous curves satisfying $\sup _{n \in \mathbb{N}} \mathscr{L}\left(\gamma_{n}\right)<\infty$ we obtain (21) in case $j p \geq 2$.

The first two parts of this result, whose second part bases on [He97, Lemma 1.4.1], were taken from [Rt04, Lemma 1.17]. The claim is also true for non-injective curves.

Proof. (i) For $\left[a^{\prime}, b^{\prime}\right] \subset X$ let $k \in \mathbb{N}$, and $a^{\prime} \leq t_{0}<\cdots<t_{k} \leq b^{\prime}$ be given. By definition we have

$$
\begin{aligned}
\mathscr{L}\left(\left.\gamma_{n}\right|_{\left[a^{\prime}, b^{\prime}\right]}\right) & \geq \sum_{i=1}^{k}\left|\gamma_{n}\left(t_{i}\right)-\gamma_{n}\left(t_{i-1}\right)\right| \\
& \geq \sum_{i=1}^{k}\left|\gamma_{0}\left(t_{i}\right)-\gamma_{0}\left(t_{i-1}\right)\right|-\left|\gamma_{n}\left(t_{i}\right)-\gamma_{0}\left(t_{i}\right)\right|-\left|\gamma_{n}\left(t_{i-1}\right)-\gamma_{0}\left(t_{i-1}\right)\right|,
\end{aligned}
$$

and this yields

$$
\liminf _{n \rightarrow \infty} \mathscr{L}\left(\left.\gamma_{n}\right|_{\left[a^{\prime}, b^{\prime}\right]}\right) \geq \sum_{i=1}^{k}\left|\gamma_{0}\left(t_{i}\right)-\gamma_{0}\left(t_{i-1}\right)\right| .
$$

Taking the supremum of decompositions $a^{\prime} \leq t_{0}<\cdots<t_{k} \leq b^{\prime}$ and $k \in \mathbb{N}$, we arrive at (20).
(ii) Passing to a subsequence we may assume the existence of $\lim _{n \rightarrow \infty} E^{j, p}\left(\gamma_{n}\right) \in[0, \infty]$. If $E^{j, p}\left(\gamma_{n}\right) \rightarrow \infty$ nothing is to prove. So let $\lim _{n \rightarrow \infty} E^{j, p}\left(\gamma_{n}\right)$ be finite. This implies $E^{j, p}\left(\gamma_{n}\right)<\infty$ for almost every $n \in \mathbb{N}$. Now we obtain $\left|\gamma_{n}(s)-\gamma_{n}(t)\right| \rightarrow\left|\gamma_{0}(s)-\gamma_{0}(t)\right|$ pointwise for all $s, t \in[a, b]$ such that

$$
\frac{1}{\left|\gamma_{0}(s)-\gamma_{0}(t)\right|^{j}}=\lim _{n \rightarrow \infty} \frac{1}{\left|\gamma_{n}(s)-\gamma_{n}(t)\right|^{j}} .
$$

On the other hand we infer $D_{\gamma_{0}}(s, t) \leq \liminf _{n \rightarrow \infty} D_{\gamma_{n}}(s, t)$ pointwise for all $s, t \in[a, b]$ from (20), i. e.

$$
\frac{1}{D_{\gamma_{0}}(s, t)^{j}} \geq \quad \liminf _{n \rightarrow \infty} \frac{1}{D_{\gamma_{n}}(s, t)^{j}} .
$$

For almost every $s \in[a, b]$ and $|h| \ll 1$ we arrive at

$$
\begin{aligned}
v_{0}=\left|\dot{\gamma}_{0}(s)\right| \stackrel{h \rightarrow 0}{\longleftrightarrow} & \left|\frac{\gamma_{0}(s+h)-\gamma_{0}(s)}{h}\right| \\
& \leq \lim _{n \rightarrow \infty}\left|\frac{\gamma_{n}(s+h)-\gamma_{n}(s)}{h}\right| \leq \liminf _{n \rightarrow \infty} v_{n}
\end{aligned}
$$

A formal calculation yields

$$
\begin{aligned}
& \left(\frac{1}{\left|\gamma_{0}(s)-\gamma_{0}(t)\right|^{j}}-\frac{1}{D_{\gamma_{0}}(s, t)^{j}}\right)^{p}\left|\dot{\gamma}_{0}(s)\right|\left|\dot{\gamma}_{0}(t)\right| \\
& \quad \leq \liminf _{n \rightarrow \infty}\left(\frac{1}{\left|\gamma_{n}(s)-\gamma_{n}(t)\right|^{j}}-\frac{1}{D_{\gamma_{n}}(s, t)^{j}}\right)^{p}\left|\dot{\gamma}_{n}(s)\right|\left|\dot{\gamma}_{n}(t)\right|
\end{aligned}
$$

for $\mathcal{L}^{2}$-almost every $s, t \in[a, b]$. By assumption the double integrals $E^{j, p}\left(\gamma_{n}\right)$ exist, and their values are finite. Fatou's lemma implies existence and finiteness of the right hand side, implying (21).
(iii) Let $\lim _{n \rightarrow \infty} E^{j, p}\left(\gamma_{n}\right)$ be finite and moreover $E^{j, p}\left(\gamma_{n}\right) \leq B<\infty$. We assume that $\left(\gamma_{n}\right)_{n \in \mathbb{N}}$ are open curves, the proof goes parallel for closed curves. Let $\tilde{\gamma}_{n}:[0,1] \rightarrow \mathbb{R}^{3}$ the reparametrization of $\gamma_{n}$ to constant velocity $\left|\dot{\tilde{\gamma}}_{n}\right|=\mathscr{L}\left(\gamma_{n}\right)$. Since the elements $\gamma_{n}$ are Lipschitz continuous with a uniform constant $\sup _{n \in \mathbb{N}} \mathscr{L}\left(\gamma_{n}\right)$, the sequence is uniformly equicontinuous. By

$$
\left|\tilde{\gamma}_{n}(t)\right| \leq \underbrace{\left|\tilde{\gamma}_{n}(t)-\tilde{\gamma}_{n}(0)\right|}_{\leq \sup _{n \in \mathbb{N}} \mathscr{L}\left(\gamma_{n}\right)}+\underbrace{\left|\tilde{\gamma}_{n}(0)-\tilde{\gamma}_{0}(0)\right|}_{\rightarrow 0 \text { as } n \rightarrow \infty}+\left|\tilde{\gamma}_{0}(0)\right|
$$

it is also uniformly bounded. So, by the Arzelà-Ascoli theorem, we may pass to a subsequence converging to some $\tilde{\gamma} \in C^{0}\left([0,1], \mathbb{R}^{3}\right)$ in the $C^{0}$-topology. The fact that the $\gamma_{n}$ are also bi-Lipschitz continuous with uniform constant max $\left(B, \sup _{n \in \mathbb{N}} \mathscr{L}\left(\gamma_{n}\right)\right)$ (see Lemma 14) is transferred to $\tilde{\gamma}$, which is in particular regular absolutely continuous. Using (ii) we arrive at $E^{j, p}(\tilde{\gamma}) \leq \liminf _{n \rightarrow \infty} E^{j, p}\left(\tilde{\gamma}_{n}\right)=\liminf _{n \rightarrow \infty} E^{j, p}\left(\gamma_{n}\right)$. We conclude this proof by verifying $E^{j, p}\left(\gamma_{0}\right)=E^{j, p}(\tilde{\gamma})$. It suffices to show Image $\gamma_{0}=$ Image $\tilde{\gamma}$. For Image $\gamma_{0} \subset$ Image $\tilde{\gamma}$ consider $y \in$ Image $\gamma_{0}$, so there is a $t \in[0,1]$ satisfying $y=\gamma_{0}(t)=\lim _{n \rightarrow \infty} \gamma_{n}(t)$. Furthermore there is a sequence $\left(t_{n}\right)_{n \in \mathbb{N}} \subset[0,1]$ with $\gamma_{n}(t)=\tilde{\gamma}_{n}\left(t_{n}\right)$. By compactness there is some $t_{0} \in[0,1]$ such that passing to a subsequence we obtain $t_{n} \rightarrow t_{0}$. Now

$$
\left|\gamma_{0}(t)-\tilde{\gamma}\left(t_{0}\right)\right| \leq \underbrace{\left|\gamma_{n}(t)-\gamma_{0}(t)\right|}_{\rightarrow 0 \text { as } n \rightarrow \infty}+\underbrace{\left|\tilde{\gamma}_{n}\left(t_{0}\right)-\tilde{\gamma}\left(t_{0}\right)\right|}_{\rightarrow 0 \text { as } n \rightarrow \infty}+\underbrace{\left|\gamma_{n}(t)-\tilde{\gamma}_{n}\left(t_{0}\right)\right|}_{\substack{\leq\left|\tilde{\gamma}_{n}\left(t_{n}\right)-\tilde{\gamma}_{n}\left(t_{0}\right)\right| \\ \leq \text { const }\left|t_{n}-t_{0}\right| \\ \rightarrow 0 \text { as } n \rightarrow \infty}}
$$

For the other direction Image $\gamma_{0} \supset$ Image $\tilde{\gamma}$ we use the fact that $\tilde{\gamma}$ is injective due to biLIPSCHITZ continuity, so $\tilde{\gamma}^{-1} \circ \gamma_{0}:[a, b] \rightarrow[0,1]$ is well-defined. Since $\tilde{\gamma}$ is continuous on a compact interval, its inverse $\tilde{\gamma}^{-1}$ is also continuous, hence $\tilde{\gamma}^{-1} \circ \gamma_{0}$ is continuous. Because of

$$
\gamma_{0}(a)=\lim _{n \rightarrow \infty} \gamma_{n}(a)=\lim _{n \rightarrow \infty} \tilde{\gamma}_{n}(0)=\tilde{\gamma}(0)
$$

we arrive at $\left(\tilde{\gamma}^{-1} \circ \gamma_{0}\right)(a)=0$ and by the same way at $\left(\tilde{\gamma}^{-1} \circ \gamma_{0}\right)(b)=1$. Using the mean value theorem we obtain for any $y=\tilde{\gamma}(t), t \in[0,1]$, a $\tau \in[a, b]$ satisfying $\left(\tilde{\gamma}^{-1} \circ \gamma_{0}\right)(\tau)=t$, hence $\gamma_{0}(\tau)=\tilde{\gamma}(t)$.

## $j p>2$

For the case $j p>2$ O'Hara showed that finite-energy $C^{1}$-curves are in fact $C^{1, \alpha / 2}$ [O'H94, Thm. 1.11]. We will prove that the same is true also for $C^{0,1}$-curves. By our framework used for arriving at pointwise differentiability we obtain the step " $C^{1} \Rightarrow$ $C^{1, \alpha / 2 "}$ almost immediately, without carrying out a geometric argument as the treatment of "solid cylinders" conducted in O'Hara's proof [O'H94, Sublemma 1.10], cf. [O'H03, pp. $67-74]$ for a more detailed version.

Again we will provide results for both closed and open curves. Let $X_{1}:=\mathbb{R} / \mathbb{Z}$ or $[0,1]$ respectively. Proving the second part of the Main Theorem we start with some technical preliminaries.

Lemma 16 ([Sem91, Lemma 8.5 revised]). For $\ell>0$ let $a:[0, \ell] \rightarrow \mathbb{R}^{n}$ be a curve parametrized by arc-length, where $P:=a(0), Q:=a(\ell)$. Then we obtain for all $t \in[0, \ell]$

$$
\left|a(t)-\left(P+\frac{t}{\ell}(Q-P)\right)\right| \leq 3 \ell\left(\frac{\ell-|P-Q|}{\ell}\right)^{1 / 2}
$$

The norm on the left hand side cannot be estimated to a power of $\ell-|P-Q|$ greater than $\frac{1}{2}$, for e. g. the arc-length parametrization of $a_{h}:[0,2] \rightarrow \mathbb{R}^{2}, t \mapsto h(t, \min (t, 2-t))$, satisfies


$$
\frac{\left|a_{h}(1)-\left(P+\frac{1}{2}(Q-P)\right)\right|}{(\ell-|P-Q|)^{\beta}}=h^{1-\beta}\left(\frac{h}{2 \sqrt{1+h^{2}}-2}\right)^{\beta}=h^{1-\beta}\left(\frac{\sqrt{1+h^{2}}+1}{2 h}\right)^{\beta} \geq h^{1-2 \beta}
$$

which tends to infinity as $h \searrow 0$ if $\beta>\frac{1}{2}$.
Proof. Applying a rotation and a translation we may assume $P=0, Q=|P-Q| e_{n}$. For $t \in[0, \ell]$ we find the following estimate for the vector $\hat{a}(t):=\left(a_{1}(t), \ldots, a_{n-1}(t)\right) \in$ $\mathbb{R}^{n-1}$.

$$
\begin{aligned}
&|\hat{a}(t)| \leq \int_{0}^{\ell}\left|\left(\dot{a}_{1}(t), \ldots, \dot{a}_{n-1}(t)\right)\right| \mathrm{d} t \leq \sqrt{\ell}\left(\int_{0}^{\ell}\left|\left(\dot{a}_{1}(t), \ldots, \dot{a}_{n-1}(t)\right)\right|^{2} \mathrm{~d} t\right)^{1 / 2} \\
& \stackrel{|\dot{a}|=1}{=} \sqrt{\ell}\left(\int_{0}^{\ell}\left(1-\dot{a}_{n}(t)^{2}\right) \mathrm{d} t\right)^{1 / 2} \leq \sqrt{\ell}\left(2 \int_{0}^{\ell}\left(1-\dot{a}_{n}(t)\right) \mathrm{d} t\right)^{1 / 2} \\
&=\sqrt{2 \ell}(\ell-|P-Q|)^{1 / 2} \leq \sqrt{2} \ell\left(\frac{\ell-|P-Q|}{\ell}\right)^{1 / 2} .
\end{aligned}
$$

Now $\left|a_{n}(\ell)\right|-a_{n}(t) \leq\left|a_{n}(\ell)-a_{n}(t)\right| \leq \ell-t$ yields $a_{n}(t) \geq|P-Q|-(\ell-t)$, which leads to

$$
a_{n}(t)-\frac{t}{\ell}|P-Q| \geq(\ell-|P-Q|)\left(\frac{t}{\ell}-1\right) \geq-(\ell-|P-Q|) .
$$

On the other hand, $a_{n}(t) \leq|a(t)| \leq t$ implies

$$
a_{n}(t)-\frac{t}{\ell}|P-Q| \leq t-\frac{t}{\ell}|P-Q|=\frac{t}{\ell}(\ell-|P-Q|) \leq \ell-|P-Q|,
$$

hence $\left|a_{n}(t)-\frac{t}{\ell}\right| P-Q| | \leq \ell\left(\frac{\ell-|P-Q|}{\ell}\right)$. Using the estimate for $\hat{a}(t)$ we conclude

$$
\left|a(t)-\frac{t}{\ell}\right| P-Q| | \leq \ell\left(\frac{\ell-|P-Q|}{\ell}\right)+\sqrt{2} \ell\left(\frac{\ell-|P-Q|}{\ell}\right)^{1 / 2} .
$$

By $x \leq \sqrt{x}$ for $x \in[0,1]$ we obtain the result.

Lemma 17. Let $\gamma \in C^{0,1}\left(X_{1}, \mathbb{R}^{3}\right)$ be parametrized by arc-length. Suppose that there are numbers $\alpha>0, \varrho_{0} \in\left(0, \frac{1}{2}\right], C<\infty$ such that for any $\varrho \in\left(0, \varrho_{0}\right]$ the quantity

$$
\kappa(\varrho) \quad:=\sup \left\{\frac{|s-t|}{|\gamma(s)-\gamma(t)|}-1\left|s, t \in X_{1},|s-t| \leq \varrho\right\}\right.
$$

fulfills the estimate

$$
\begin{equation*}
\kappa(\varrho) \leq C \varrho^{\alpha} . \tag{22}
\end{equation*}
$$

Then there is an $\varepsilon_{0}=\varepsilon_{0}\left(\alpha, \varrho_{0}, C\right)>0$, such that all $x, z, \xi, \zeta \in X_{1}$ with $|x-z| \leq \varepsilon_{0}$, $x \leq \xi<\zeta \leq z$, and $|\xi-\zeta| \geq \frac{1}{2}|x-z|$ satisfy

$$
\left|\frac{\gamma(z)-\gamma(x)}{|\gamma(z)-\gamma(x)|}-\frac{\gamma(\zeta)-\gamma(\xi)}{|\gamma(\zeta)-\gamma(\xi)|}\right| \leq 48 \sqrt{C} \cdot|x-z|^{\alpha / 2} .
$$

Note that the restriction $|s-t| \leq \varrho_{0} \leq \frac{1}{2}$ implies $|s-t|=D_{\gamma}(s, t)$ and

$$
\begin{equation*}
|s-t| \leq(1+\kappa(|s-t|)) \cdot|\gamma(s)-\gamma(t)| \quad \text { for all }|s-t| \leq \varrho_{0} . \tag{23}
\end{equation*}
$$

Proof. We choose $\varepsilon_{0} \in\left(0, \varrho_{0}\right]$ so small that

$$
\begin{equation*}
\kappa\left(\varepsilon_{0}\right) \leq 1 \tag{24}
\end{equation*}
$$

For $x, z, \xi, \zeta \in X_{1}$ as in the assumptions we set

$$
\begin{aligned}
a & :=\gamma(\zeta)-\gamma(\xi), \\
b & :=\frac{\zeta-\xi}{z-x}[\gamma(z)-\gamma(x)] .
\end{aligned}
$$

Decomposing

$$
a-b=\left[\gamma(\zeta)-\left(\frac{\zeta-x}{z-x} \gamma(z)+\frac{z-\zeta}{z-x} \gamma(x)\right)\right]-\left[\gamma(\xi)-\left(\frac{\xi-x}{z-x} \gamma(z)+\frac{z-\xi}{z-x} \gamma(x)\right)\right]
$$

we apply Lemma 16 with $\ell=z-x, a=\gamma(\cdot-x)$, and $t=\zeta$ or $\xi$ respectively. This yields

$$
\begin{aligned}
|a-b| & \leq 6|z-x|\left(\frac{|z-x|-|\gamma(z)-\gamma(x)|}{|z-x|}\right)^{1 / 2} \stackrel{(23)}{\leq} 6|z-x|\left(\frac{|z-x|-\frac{|z-x|}{1+\kappa(|z-x|)}}{|z-x|}\right)^{1 / 2} \\
& =6|z-x|\left(1-\frac{1}{1+\kappa(|z-x|)}\right)^{1 / 2} \leq 6|z-x| \cdot \kappa(|z-x|)^{1 / 2} \\
& \stackrel{(22)}{\leq} 6 \sqrt{C}|z-x|^{\alpha / 2+1} .
\end{aligned}
$$

Since

$$
|\gamma(\xi)-\gamma(\zeta)| \stackrel{(23)}{\geq} \frac{|\xi-\zeta|}{1+\kappa(|\xi-\zeta|)} \stackrel{(24)}{\geq} \frac{1}{2}|\xi-\zeta| \quad \frac{1}{4}|z-x|
$$

we obtain recalling $\left|\frac{a}{|a|}-\frac{b}{|b|}\right| \leq 2 \frac{|a-b|}{|a|}$

$$
\left|\frac{\gamma(z)-\gamma(x)}{|\gamma(z)-\gamma(x)|}-\frac{\gamma(\zeta)-\gamma(\xi)}{|\gamma(\zeta)-\gamma(\xi)|}\right| \leq 2 \frac{|a-b|}{|a|} \leq 48 \sqrt{C}|z-x|^{\alpha / 2} .
$$

Lemma 18. The hypotheses of Lemma 17 imply $\gamma \in C^{1, \alpha / 2}\left(X_{1}, \mathbb{R}^{3}\right)$.

Proof. For $x, z \in X_{1}$ with $x<z \leq x+\varepsilon_{0}$ and $x \leq s<t \leq z$ let $k \in \mathbb{N}$ be such that

$$
2^{-k+1}|z-x| \geq|t-s|>2^{-k}|z-x| .
$$

Then there are $\xi_{l}, \zeta_{l} \in[x, z], \xi_{l}<\zeta_{l}, l=0, \ldots, k$ satisfying $\left[\xi_{0}, \zeta_{0}\right]=[x, z],\left[\xi_{k}, \zeta_{k}\right]=$ $[s, t],\left[\xi_{l}, \zeta_{l}\right] \subset\left[\xi_{l-1}, \zeta_{l-1}\right]$, and $\left|\zeta_{l}-\xi_{l}\right|=\frac{1}{2}\left|\zeta_{l-1}-\xi_{l-1}\right|$ for $l=0, \ldots, k-1$ and $\left|\zeta_{k}-\xi_{k}\right| \geq$ $\frac{1}{2}\left|\zeta_{k-1}-\xi_{k-1}\right|$. Applying Lemma 17 to

$$
v_{l}:=\frac{\gamma\left(\zeta_{l}\right)-\gamma\left(\xi_{l}\right)}{\left|\gamma\left(\zeta_{l}\right)-\gamma\left(\xi_{l}\right)\right|}
$$

we arrive at

$$
\begin{equation*}
\left|v_{l-1}-v_{l}\right| \leq 48 \sqrt{C} \cdot 2^{(-l+1) \alpha / 2} \cdot|z-x|^{\alpha / 2} \tag{25}
\end{equation*}
$$

for $l=1, \ldots k$. We now compute

$$
\begin{align*}
& \left|\frac{\gamma(z)-\gamma(x)}{z-x}-\frac{\gamma(t)-\gamma(s)}{t-s}\right| \\
& \leq\left|\frac{\gamma(z)-\gamma(x)}{z-x}-\frac{\gamma(z)-\gamma(x)}{|\gamma(z)-\gamma(x)|}\right|+\left|\frac{\gamma(z)-\gamma(x)}{|\gamma(z)-\gamma(x)|}-\frac{\gamma(t)-\gamma(s)}{|\gamma(t)-\gamma(s)|}\right| \\
& +\left|\frac{\gamma(t)-\gamma(s)}{|\gamma(t)-\gamma(s)|}-\frac{\gamma(t)-\gamma(s)}{t-s}\right| \\
& \stackrel{(25)}{\leq} \frac{\frac{|z-x|}{|\gamma(z)-\gamma(x)|}-1}{\frac{|z-x|}{|\gamma(z)-\gamma(x)|}}+\sum_{l=1}^{k}\left|v_{l-1}-v_{l}\right|+\frac{\frac{|t-s|}{|\gamma(t)-\gamma(s)|}-1}{\frac{|t-s|}{|\gamma(t)-\gamma(s)|}} \\
& \stackrel{(23)}{\leq} \kappa(|z-x|)+48 \sqrt{C} \cdot|z-x|^{\alpha / 2} \sum_{l=0}^{\infty}\left(2^{-\alpha / 2}\right)^{l}+\kappa(|t-s|) \\
& \stackrel{(22)}{\leq} 2 C \cdot|z-x|^{\alpha}+48 \sqrt{C} \cdot|z-x|^{\alpha / 2} \cdot \frac{1}{1-2^{-\alpha / 2}} \text {. } \\
& \leq C_{\alpha}|z-x|^{\alpha / 2}, \tag{26}
\end{align*}
$$

where $C_{\alpha}$ is a constant depending only on $\alpha$ and $C$.
Now let $y \in \mathbb{R}$ or $y \in(0,1)$ respectively and choose $\delta \in\left(0, \varepsilon_{0}\right]$ so small that $y \pm \delta \in X_{1}$. Using the last inequality with $x=y-\delta$ and $z=y+\delta$ we obtain

$$
\left|\frac{\gamma(t)-\gamma(s)}{t-s}-\frac{\gamma(x+\delta)-\gamma(x-\delta)}{2 \delta}\right| \leq C_{\alpha}(2 \delta)^{\alpha / 2}
$$

for all $y-\delta \leq s<t \leq y+\delta$. Thus $\gamma$ is differentiable in $y$.
Finally we obtain $\gamma \in C^{1, \alpha / 2}$ by

$$
\begin{aligned}
\mid \dot{\gamma}(z) & -\dot{\gamma}(x) \mid \\
& \leq\left|\frac{\gamma(z)-\gamma(x)}{z-x}-\dot{\gamma}(z)\right|+\left|\frac{\gamma(z)-\gamma(x)}{z-x}-\dot{\gamma}(x)\right| \\
& \leq \quad \lim _{h \searrow 0}\left|\frac{\gamma(z)-\gamma(x)}{z-x}-\frac{\gamma(z)-\gamma(z-h)}{h}\right|+\lim _{h \supset 0}\left|\frac{\gamma(z)-\gamma(x)}{z-x}-\frac{\gamma(x+h)-\gamma(x)}{h}\right| \\
& \leq \quad 2 C_{\alpha}|z-x|^{\alpha / 2}
\end{aligned}
$$

for all $x, z \in X_{1}$ with $|z-x| \leq \varepsilon_{0}$ and

$$
|\dot{\gamma}(z)-\dot{\gamma}(x)| \leq 2 \varepsilon_{0}^{-\alpha / 2} D_{\gamma}(z, x)^{\alpha / 2}
$$

for all $x, z \in X_{1}$ with $D_{\gamma}(z, x) \geq \varepsilon_{0}$.
Proposition 19 (Quantified bi-LiPschitz estimate for finite energy curves of unitlength, [O'H94, Prop. 1.6]). For jp $>2$ let $\gamma \in C^{0,1}\left(X_{1}, \mathbb{R}^{3}\right)$ be parametrized by arc-length with $E^{(j, p)}(\gamma) \leq B$ for some $B>0$. Then there is an $A=A(j, p, B)>0$, such that, provided $|s-t| \leq \varrho_{0}:=\min \left((2 A)^{-(p+2) /(j p-2)}, \frac{1}{2}\right)$,

$$
|s-t| \leq \frac{|\gamma(s)-\gamma(t)|}{1-A|\gamma(s)-\gamma(t)|^{(j p-2) /(p+2)}}
$$

The proof which may be found in [O'H94, pp. 49 - 51] or [O'H03, Cor. 4.2.3 (1)] also holds for open curves. The condition $|s-t| \leq \min \left((2 A)^{-(p+2) /(j p-2)}, \frac{1}{2}\right)$ guarantees that the denominator of the right hand side is bounded below by $\frac{1}{2}$ and that $|s-t|=D_{\gamma}(s, t)$ holds.

Proof of the second part of the Main Theorem. As mentioned in the introduction we may restrict ourselves to arc-length parametrized curves. Due to scaling invariance we may furthermore assume that the length of our curve is 1 . So let $\gamma \in C^{0,1}\left(X_{1}, \mathbb{R}^{3}\right)$ be parametrized by arc-length with $E^{j, p}(\gamma)<\infty$. Proposition 19 guarantees condition (22) with $C=2 A$ and $\alpha=(j p-2) /(p+2)$, for

$$
\frac{|s-t|}{|\gamma(s)-\gamma(t)|}-1 \leq \frac{A|\gamma(s)-\gamma(t)|^{(j p-2) /(p+2)}}{1-A|\gamma(s)-\gamma(t)|^{(j p-2) /(p+2)}} \leq 2 A|\gamma(s)-\gamma(t)|^{(j p-2) /(p+2)} .
$$

Now the claim follows by Lemma 18.

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    ${ }^{1}$ In general, self-repulsion does not penalize "pulling tight" of small knots, cf. [O'H92b, Thm. 3.1 (2)].

[^1]:    ${ }^{2}$ In fact, O'Hara's first version of a knot energy equals $\frac{1}{2} E-2$.
    ${ }^{3}$ Later, using the machinery of pseudo-differential operators He was able to prove $C^{\infty}$-smoothness of the $E$-minimizers [ He 00 , Cor. 5.3], thus resolving completely the regularity theory for minimizers of this particular knot energy.
    ${ }^{4}$ In [FHW94, Rem. to Cor. 1.3] the authors conjectured the existence of such a curve.

[^2]:    ${ }^{5}$ This observation also shows that the assumption $(j-2) p<1$ in Part (i) of the Main Theorem is in the case of closed curves not a restriction at all.
    ${ }^{6}$ To be more precise, on $\{(j-2) p \geq 1, j p>2\}$ we know that all curves which are at least $C^{3,1}$ (except lines, since their curvature vanishes completely) and, on the other hand, all curves which fail to be $C^{1, \alpha / 2}$ have infinite energy, see Lemma 1 (ii) and part (ii) of the Main Theorem.

[^3]:    ${ }^{7}$ The energy of a line amounts to 0 for all $j, p>0$.

[^4]:    ${ }^{8}$ Hence O'Hara's results for $j p>2$ [O'H94] automatically extend to the larger class of regular absolutely continuous curves.

